

# Probing neutrons and purported fission daughter products from gas-loaded, laser-irradiated metal-hydrogen targets

Principal Investigator: Yet-Ming Chiang, Professor, MIT

Presentation by Florian Metzler, Research Scientist, MIT

Washington, DC  
September 8, 2023



U.S. DEPARTMENT OF  
**ENERGY**

# ARPA-E LENR: Massachusetts Institute of Technology



Massachusetts Institute of Technology

## Project Title:

Probing neutrons and purported fission daughter products from gas-loaded, laser-irradiated metal-hydrogen targets

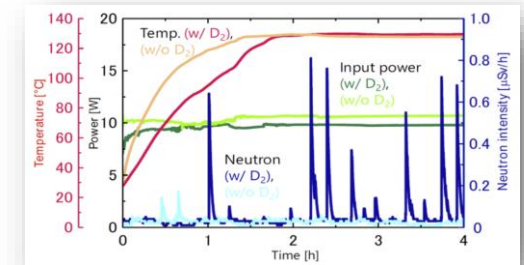
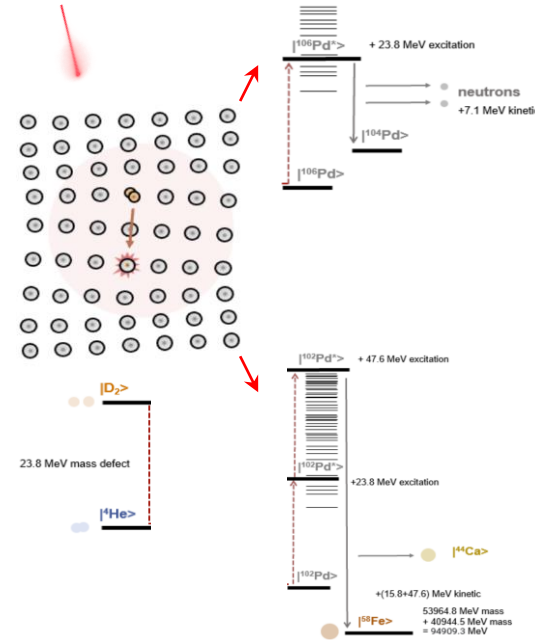
PI:

Yet-Ming Chiang  
ychiang@mit.edu

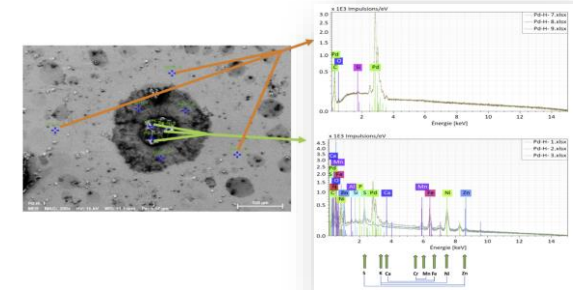
## Project Outcomes:

Unequivocal demonstration of the presence or absence of LENR phenomena through replication experiments with additional characterization and detection methods

Laser irradiation



Neutron bursts



Elemental production

Key takeaway: anomalous emitted radiation and surface isotopes have been reported in metal deuterides under low power laser irradiation and should be revisited with improved radiation detection and isotopic characterization

# Hypothesis

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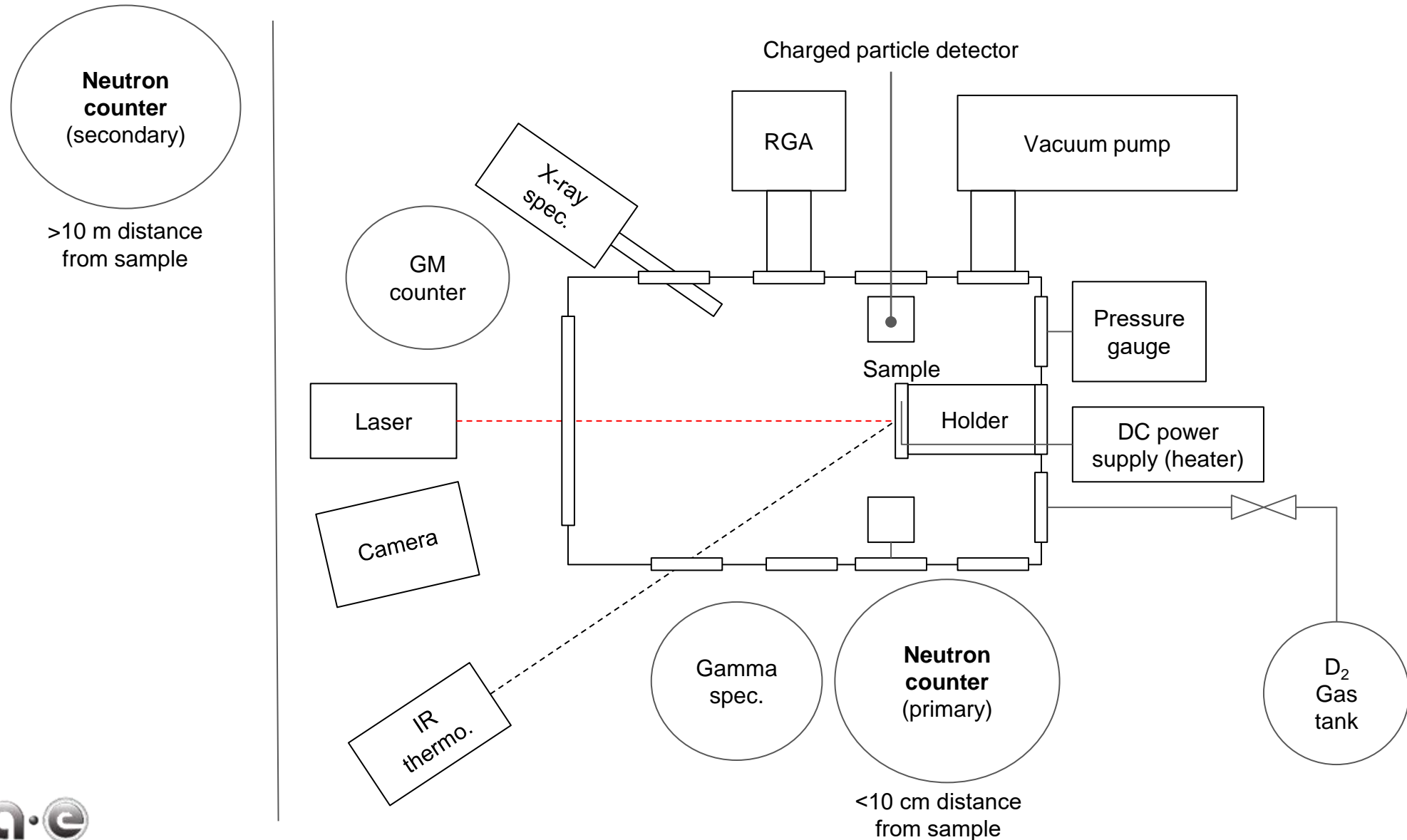
## Hypothesis:

Nuclear reactions can occur in metal deuterides at near-ambient temperature and pressure conditions under low-power laser irradiation  
(Belyukov 1991; Nassissi 1999; Mastromatteo 2016; Ushikoshi 2020; Barrowes 2022)

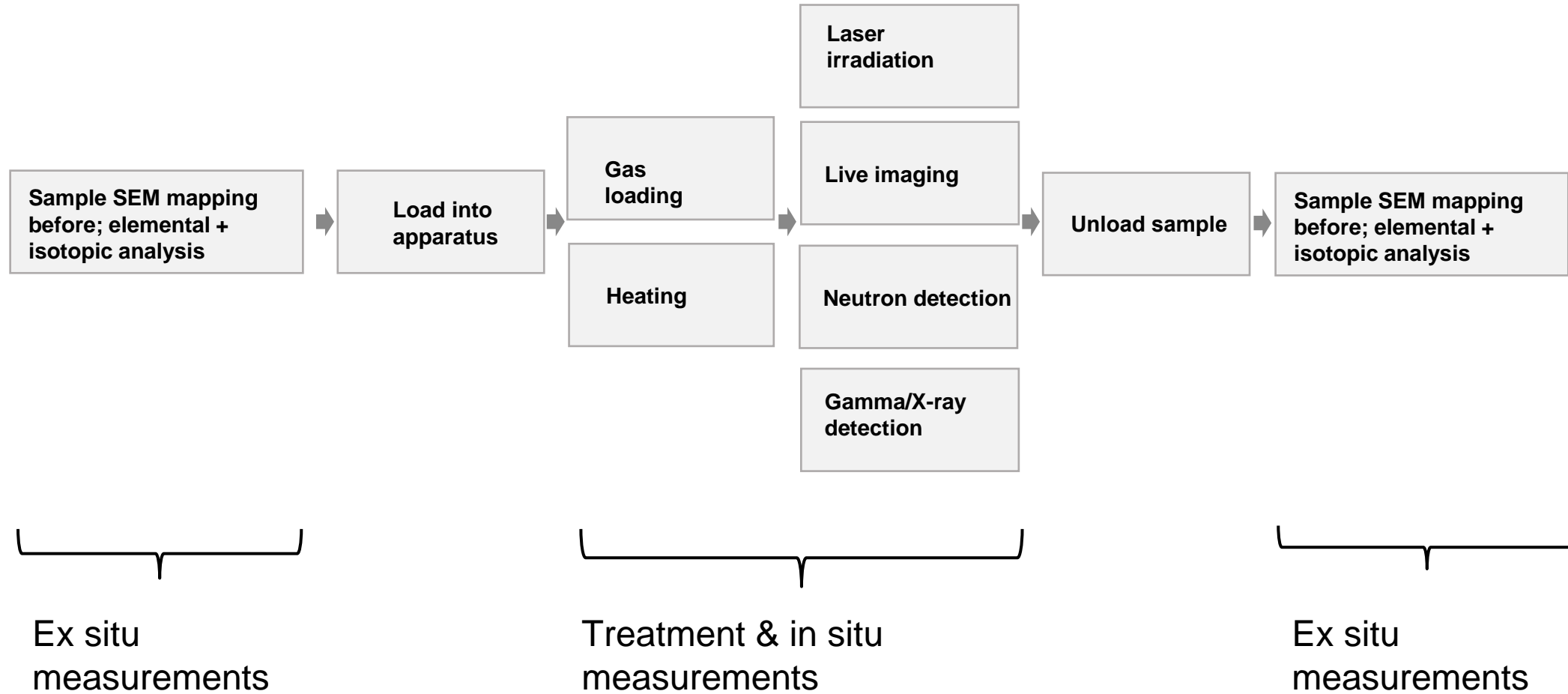
## More specifically:

- DD and HD fusion reactions can be accelerated by nonradiatively transferring corresponding transition energies (23.8 MeV and 5.5 MeV respectively) to resonant excited states of heavy lattice nuclei (e.g. Pd and Ti isotopes).
- Such transfer is enabled by shared phonon and plasmon modes that cause temporary delocalization of nuclear states in a coherence domain (delocalized nuclear excitons).
- Dicke enhancement can accelerate what are initially low transfer probabilities due to weak couplings.

# Overview of experimental setup



# Key steps



# Variables overview

## Independent variables

- **target material** (Pd, Ti);
- **target temperature** (25–175 °C)
- **chamber gas** (vacuum, H<sub>2</sub>, D<sub>2</sub>, Ar)
- **chamber pressure** (0.01–4 bar);
- **laser wavelength** (405, 594, 640, 1064 nm);
- **laser type** (pulsed or continuous-wave)
- **laser strength** (up to 1 J/pulse or 5-50 mW for CW);
- **spot size** (0.01–1 cm<sup>2</sup>).

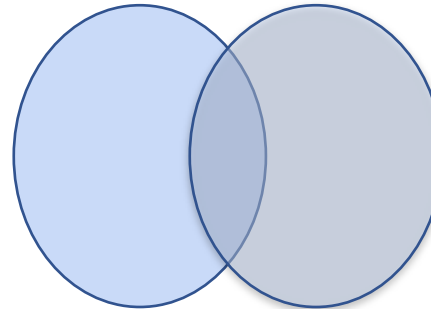
+ **potentially many hidden variables**



## Observables

Spots  
with low-Z  
elements

Bursts  
in neutron and  
gamma detectors



- **neutron** counts (He-3 counter detectors; efficiency ~20%)
- **gamma** counts >100 keV (NaI gamma spectrometers)
- **isotopic ratios** of sample surface materials (mass spectrometer + NAA; goal is accuracy  $\pm 10\%$ )

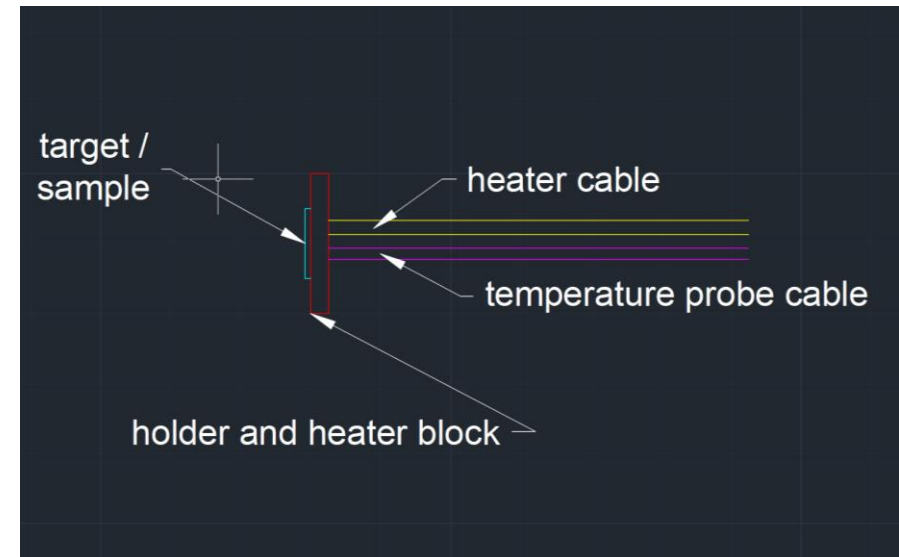
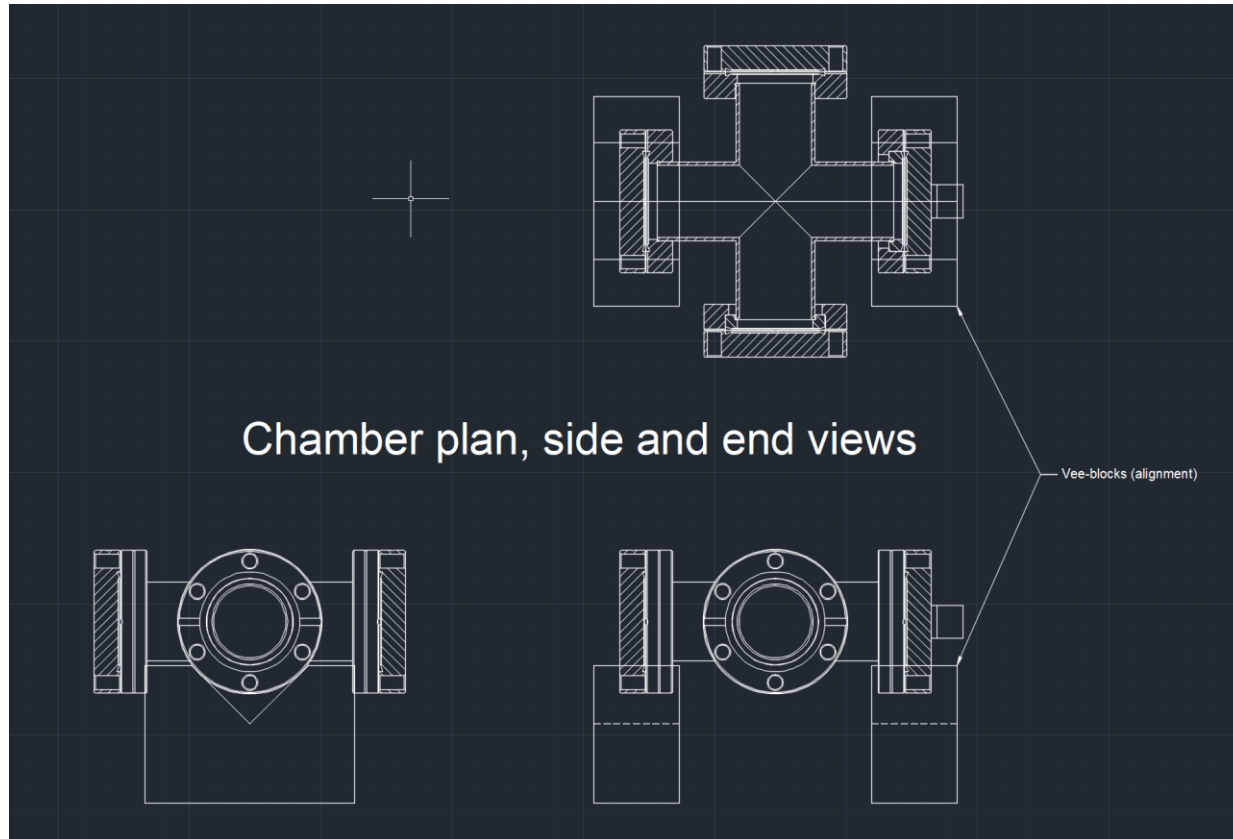


## Causes

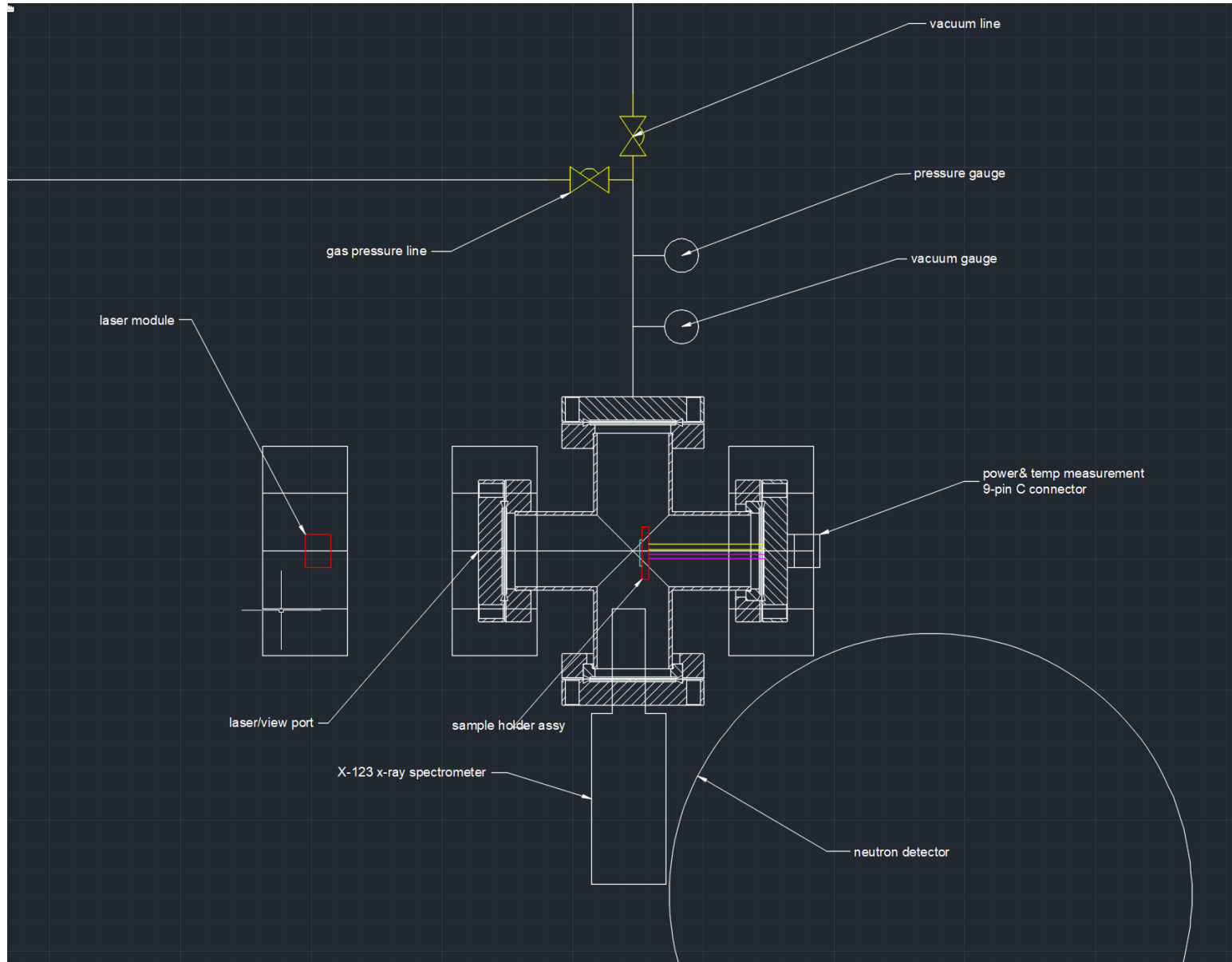
Prosaic?

Nuclear?

# Reactor design



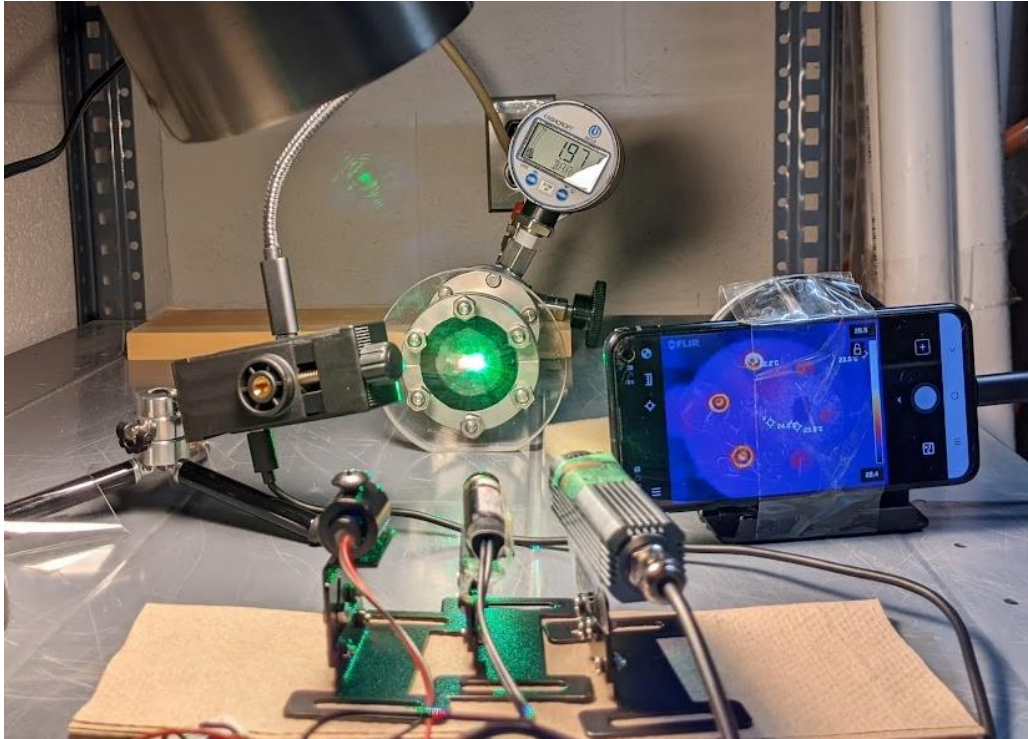
# Reactor design (cont'd)



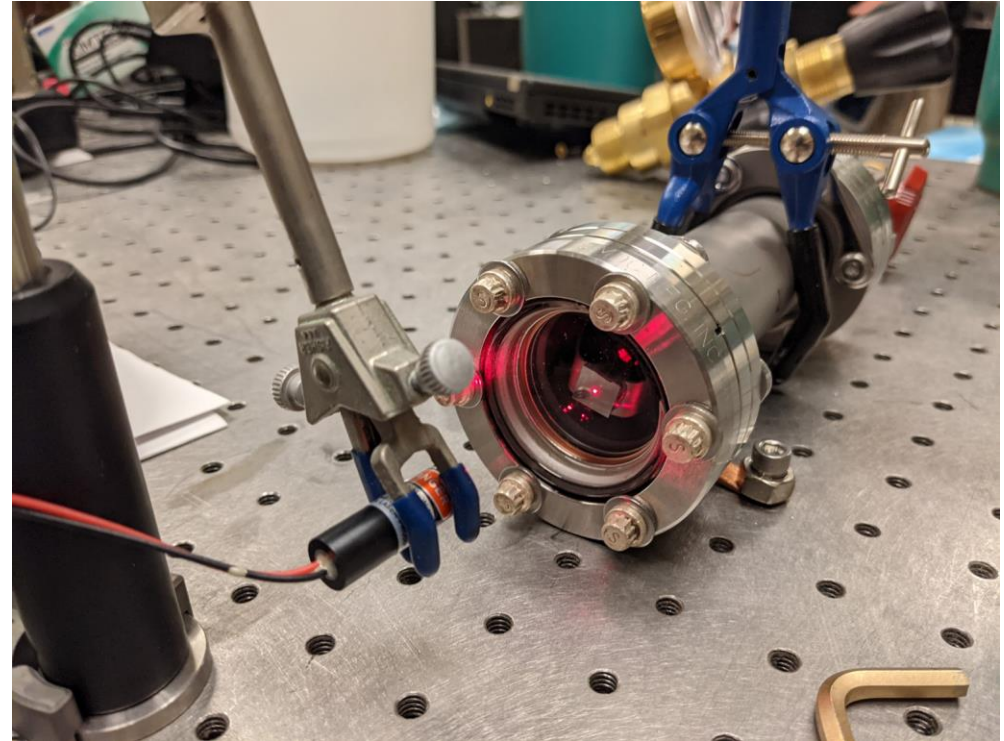


# Preliminary experiments

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At US Army EDRC



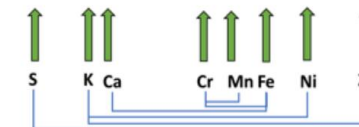
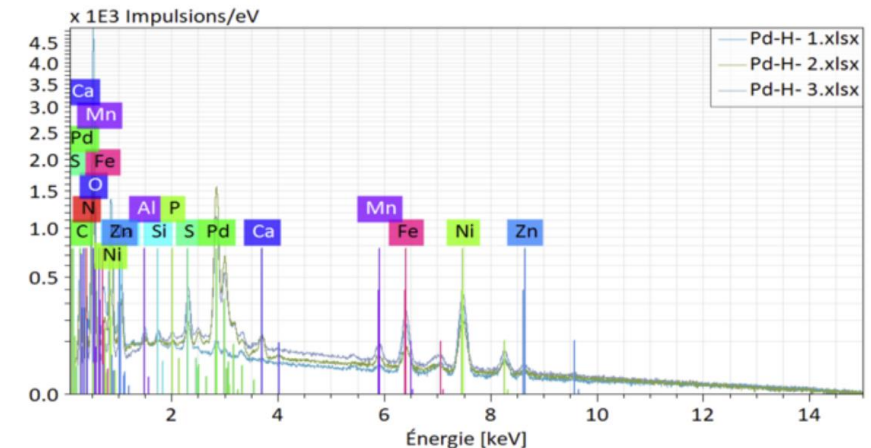
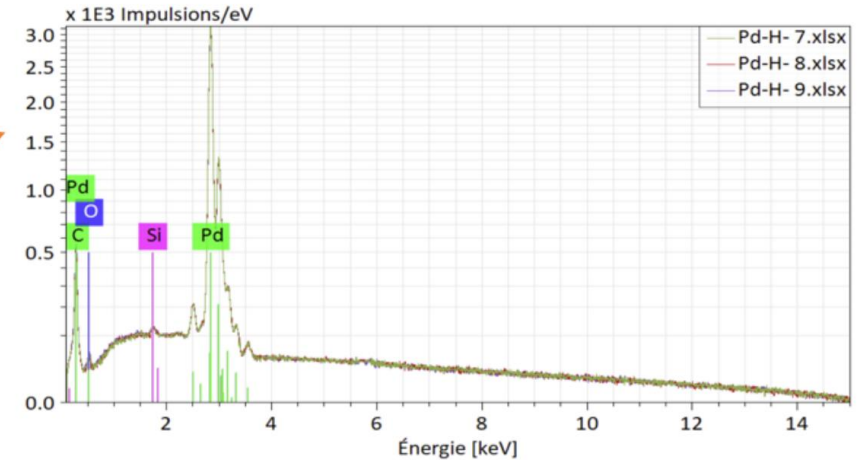
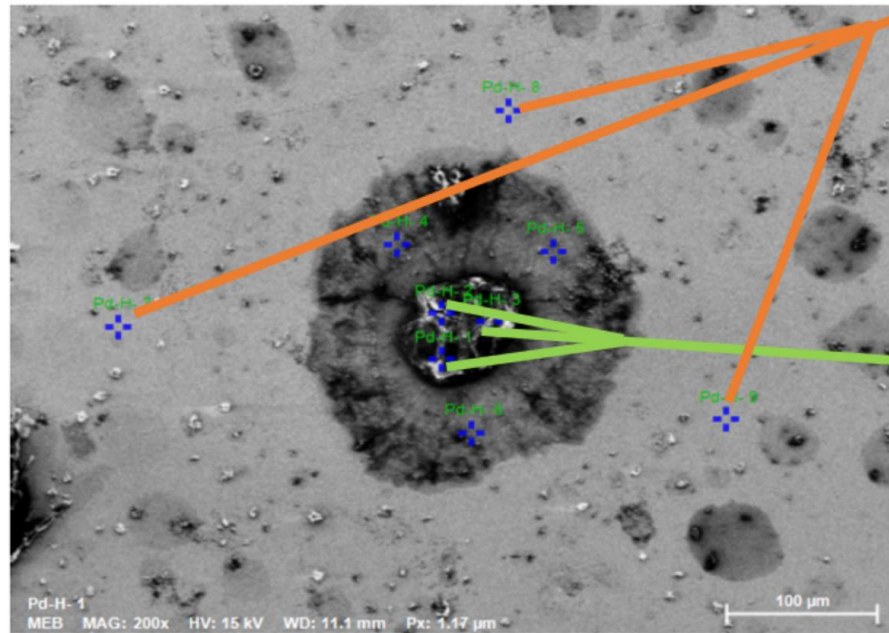
At MIT

# Data Acquisition

Measurement	Recording Method	Settings	Latency	Storage Media
SEM-EDS surface maps	Bitmaps	TBD	Before and after each experiment	Dropbox or Google Drive
Isotopic detection	Gamma spectra from Neutron activation analysis (NAA) and mass spectra from MS	TBD	Before and after each experiment	Dropbox or Google Drive
Radiation detection ( $n$ )	Neutron counts per second from $^3\text{He}$ proportional counter & liquid scintillation detector	Use of moderator to optimize sensitivity for 1-10 MeV neutrons	Continual data collection w/ time resolution of counts per second or higher	Local computer + Postgres time series database on an MIT-based server; Dropbox or Google Drive backup
Radiation detection ( $\gamma$ )	Spectra and counts >100 keV from HPGe and NaI spectrometers	Focus on 100 keV to 5 MeV range	Continual data collection w/ time resolution of counts per min or higher	Local computer + Postgres time series database on an MIT-based server; Dropbox or Google Drive backup

# Reports of surface changes in the literature

## SEM-EDS analysis of spots

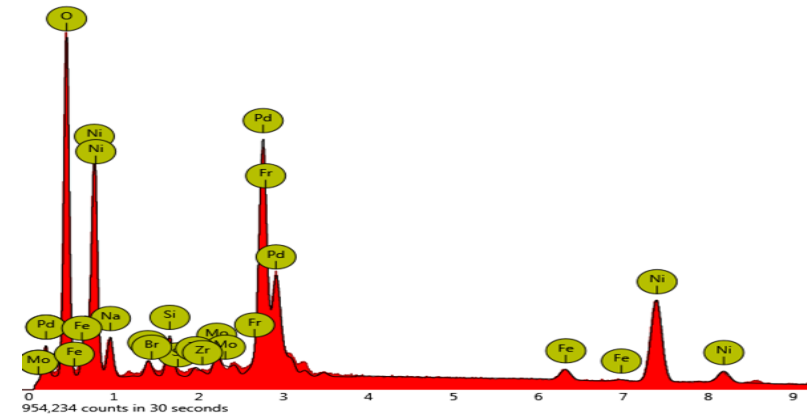
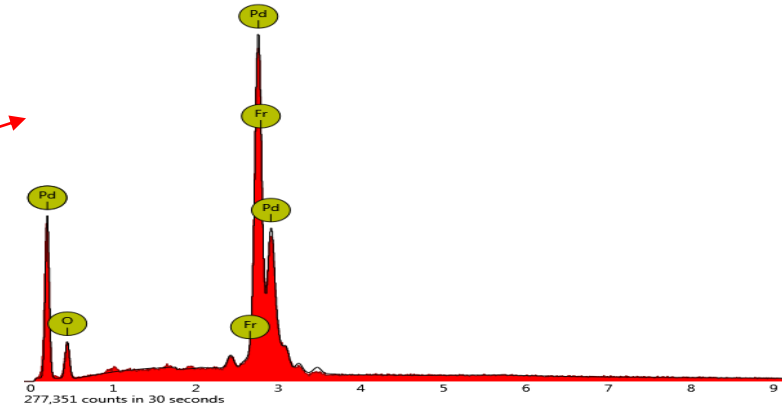
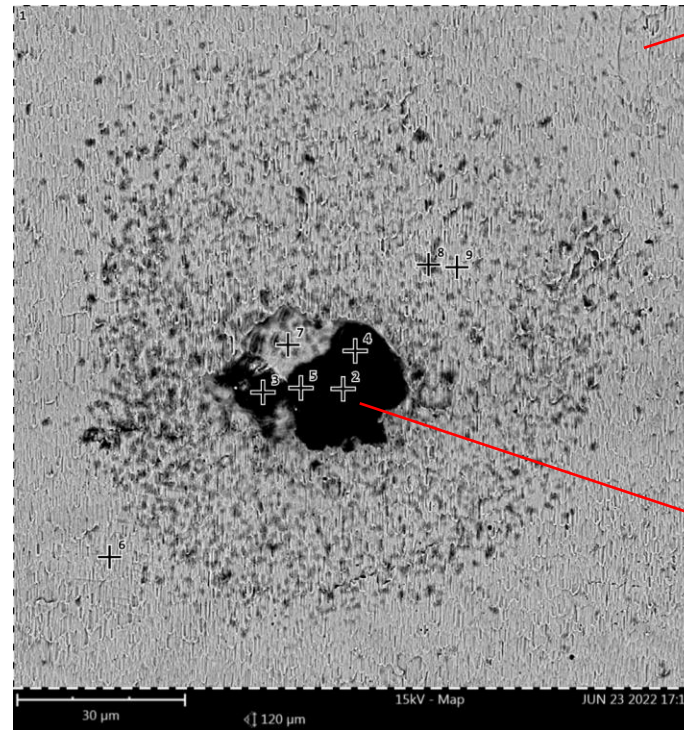


**Selection of possible stable daughter pairs from the fission of Pd isotopes:**

- Pd-102: Ni + K
- Pd-104: Zn + S
- Pd-105: Fe + Ca, Ni + K
- Pd-106: Fe + Ca, Zn + S, Fe + Cr
- Pd-108: Zn + K, Fe + Cr, Mn + Cr
- Pd-110: Fe + Cr, Mn + Mn

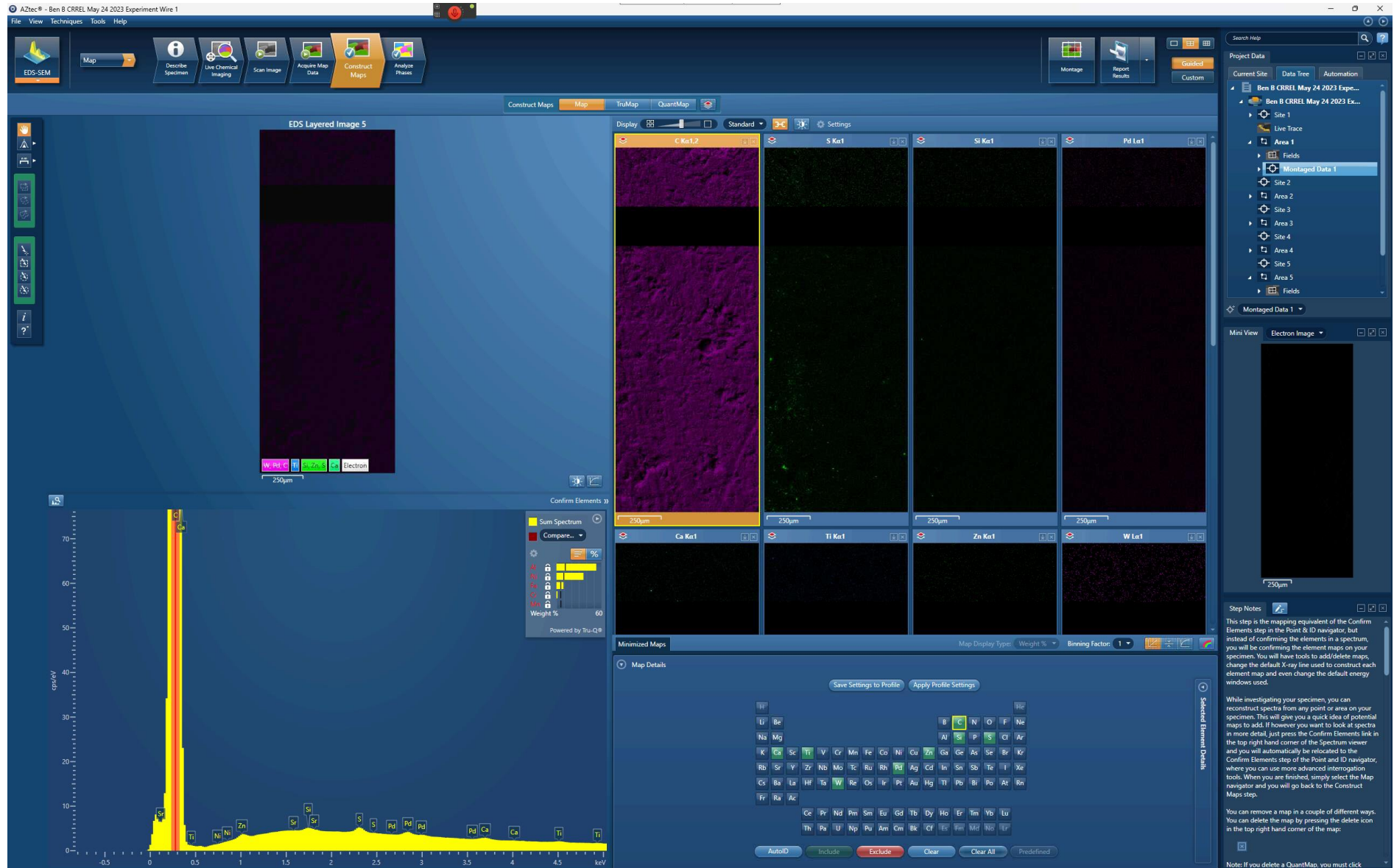
# Reports of surface changes in the literature (cont'd)

## SEM-EDS analysis of spots

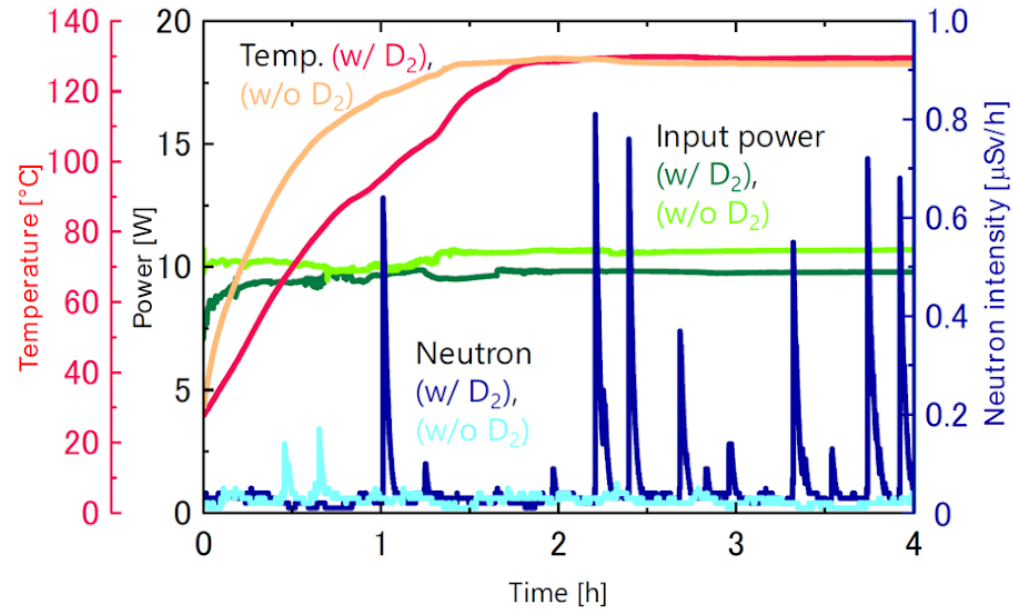


Barrowes 2022, ICCF24

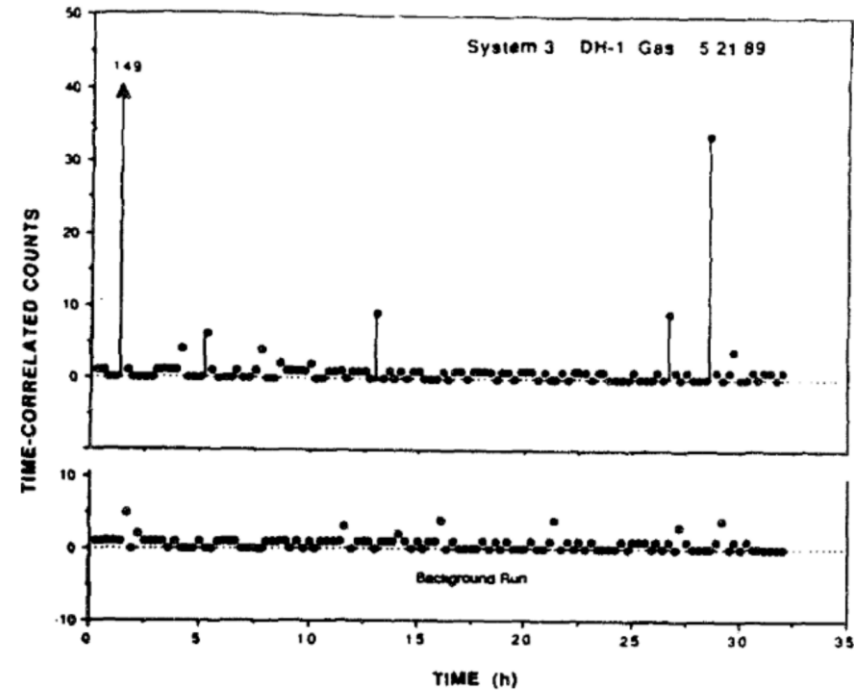
# SEM-EDS surface mapping in collaboration with Texas Tech



# Reports of neutron bursts in the literature



Uchikoshi 2020, Kyoto University



Menlove et al. 1990, Journal of Fusion Energy

# Reports of neutron bursts in the literature (cont'd)

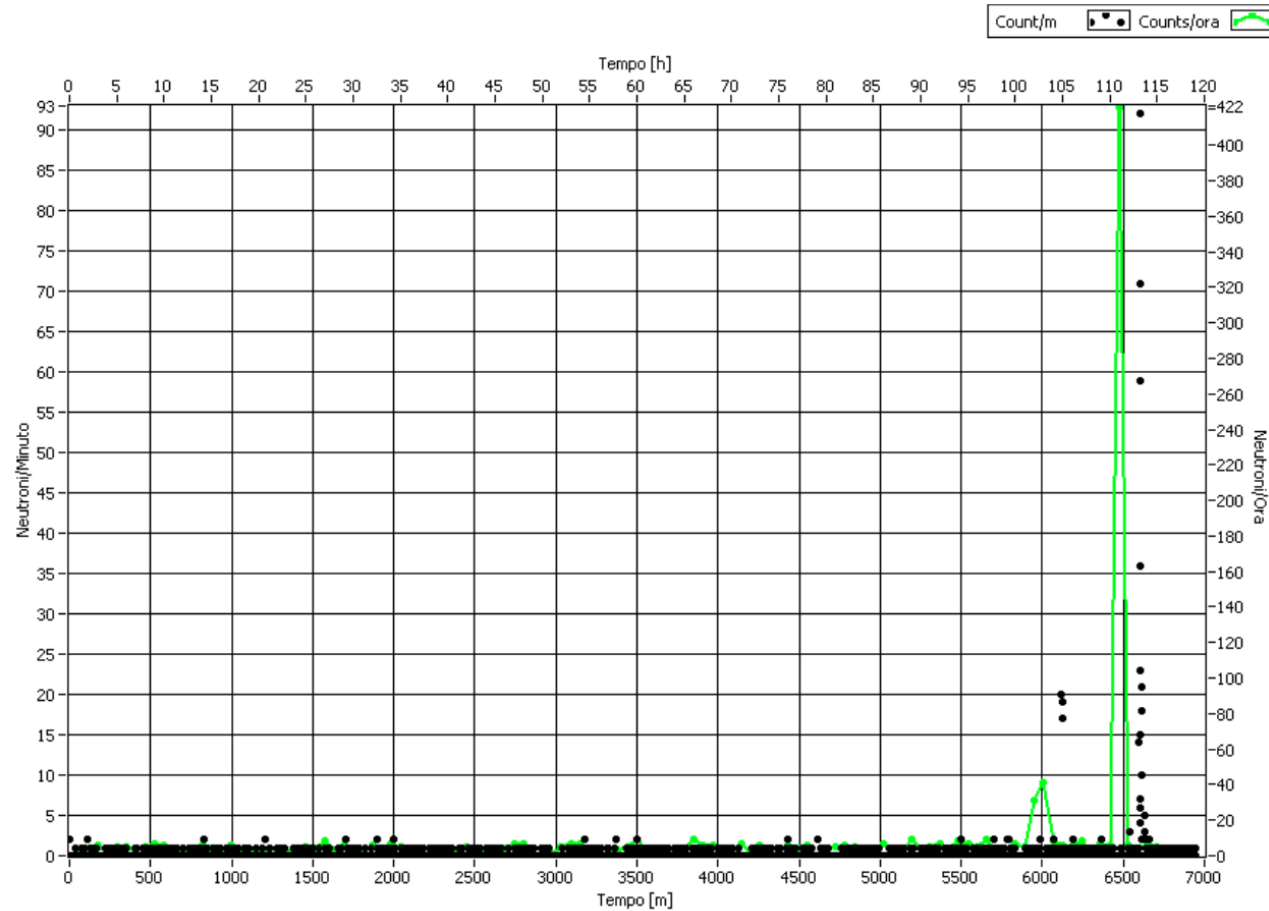
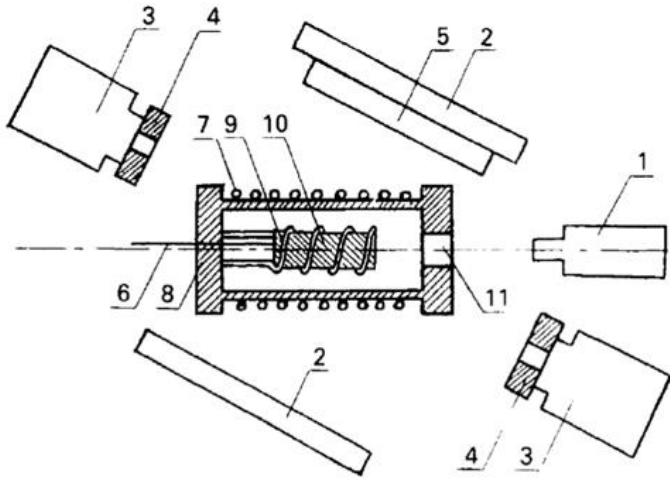


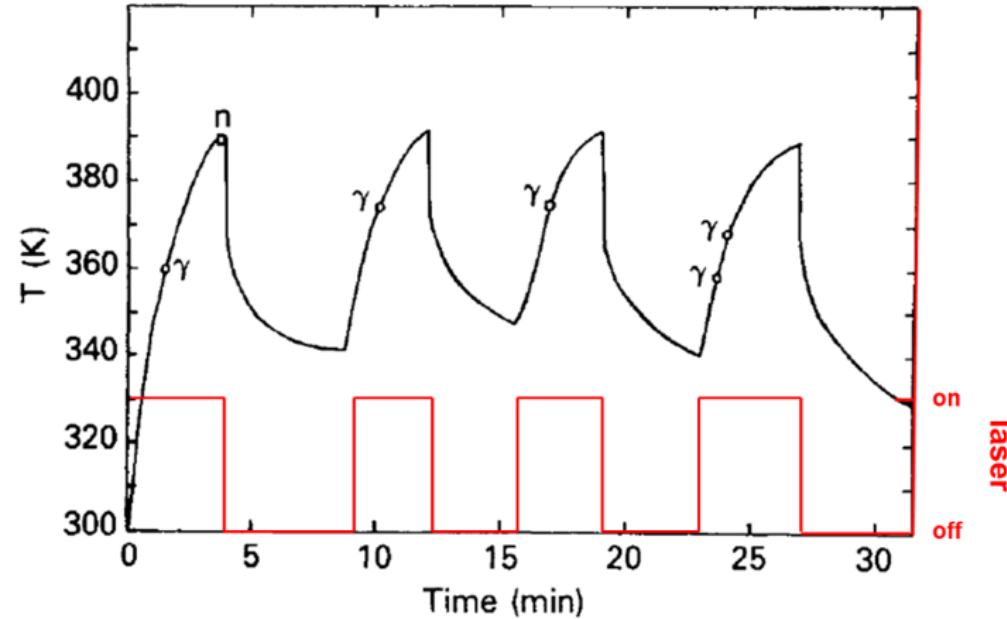
Figure 14. Screen shot of the neutron monitoring PC showing a massive neutron emission from reactor 1 after the experiment with a 405 nm laser.

Mastromatteo 2020, JCMNS

# Reports of neutron and gamma bursts correlated with irradiation



- |  |                             |
|--|-----------------------------|
| 1 Laser LTI-403                        | 6 Thermocouple              |
| 2 Neutron detector                     | 7 Electroresistance furnace |
| 3 Gamma detector                       | 8 Gas cell                  |
| 4 100-mm lead screen                   | 9 Heat exchanger            |
| 5 90-mm polyethylene neutron moderator | 10 Sample                   |
|  | 11 Quartz glass             |



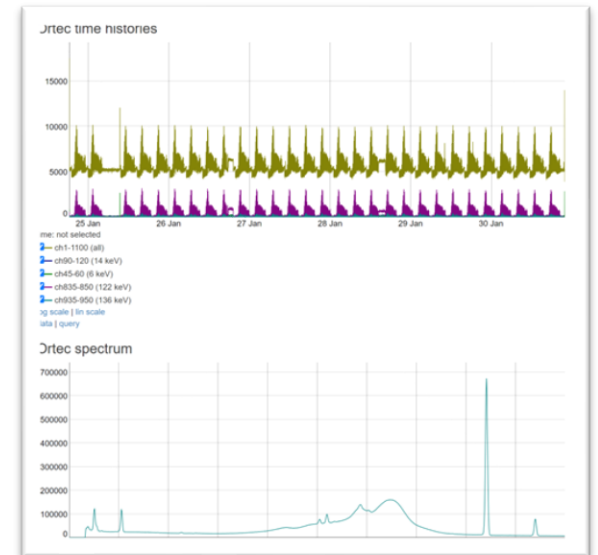
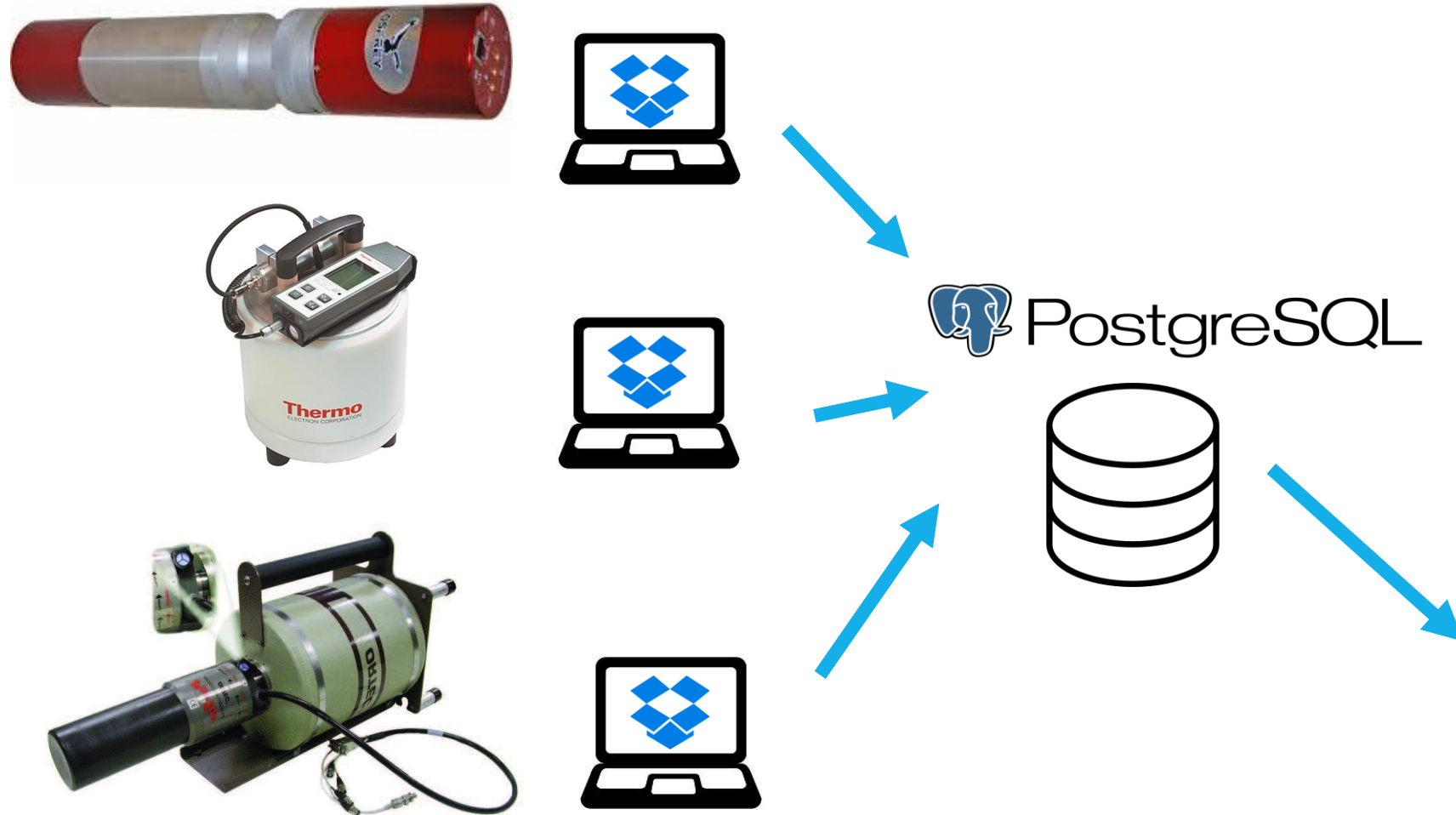
Gamma emission on the order of  $5 \times 10^3/s$   
and neutron emission on the order of  $2 \times 10^2 n/s$   
in bursts of  $<0.5 s$

(exceeding gamma background by 2x and neutron background by 18x)

Beltyukov et al. 1991, Fusion Technology



# 24/7 data streaming with universal time stamps



# Cloud-based time series optimized Postgres database

Adminer 4.7.3 4.8.1

DB: experiments  
Schema: public

SQL command Import  
Export Create table

Select: neutron1

Select data Show structure Alter table New item

Select Search Sort Limit Action

time descending  
descending

51 Select !

```
SELECT * FROM "neutron1" LIMIT 51 (0.001 s) Edit
```

<input type="checkbox"/> Modify	time	channels	id
<input type="checkbox"/> edit	2020-01-10 15:52:44.533429	{0}	180119
<input type="checkbox"/> edit	2020-01-10 15:53:07.17187	{0}	180120
<input type="checkbox"/> edit	2020-01-10 15:53:12.445233	{0}	180121
<input type="checkbox"/> edit	2020-01-10 15:53:17.704783	{0}	180122
<input type="checkbox"/> edit	2020-01-10 15:53:22.985357	{0}	180123
<input type="checkbox"/> edit	2020-01-10 15:53:28.272824	{0}	180124
<input type="checkbox"/> edit	2020-01-10 15:53:33.535181	{0}	180125
<input type="checkbox"/> edit	2020-01-10 15:53:38.80538	{0}	180126
<input type="checkbox"/> edit	2020-01-10 15:53:44.107407	{0}	180127
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<input type="checkbox"/> edit	2020-01-10 15:54:15.722133	{0}	180133
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<input type="checkbox"/> edit	2020-01-10 15:54:58.023233	{0}	180141

Page 1 2 3 4 5 ... last Whole result ~ 1,266,176 rows Selected (0) Export (~ 1,266,176)

Save Edit Clone Delete

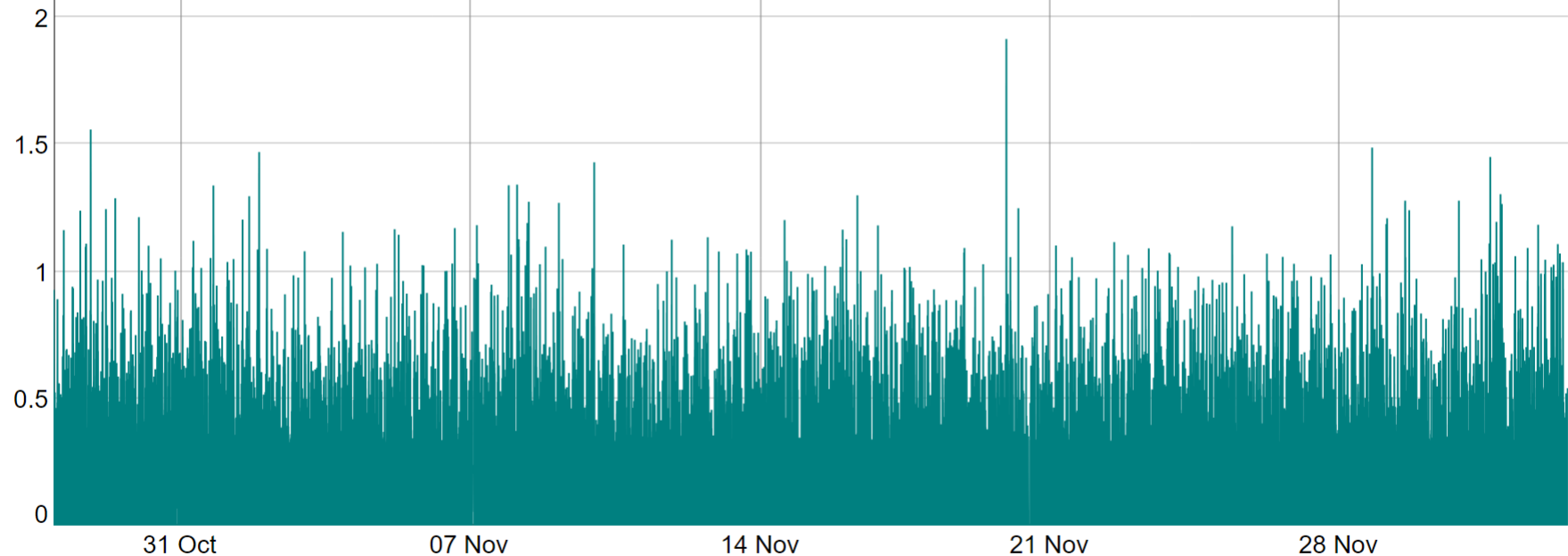


# Example: neutron background

[show editor](#) | [hide editor](#) | [clear cache and refresh](#)

## Neutrons lab

### Neutron counts



time: not selected

Neutrons

[log scale](#) | [lin scale](#)

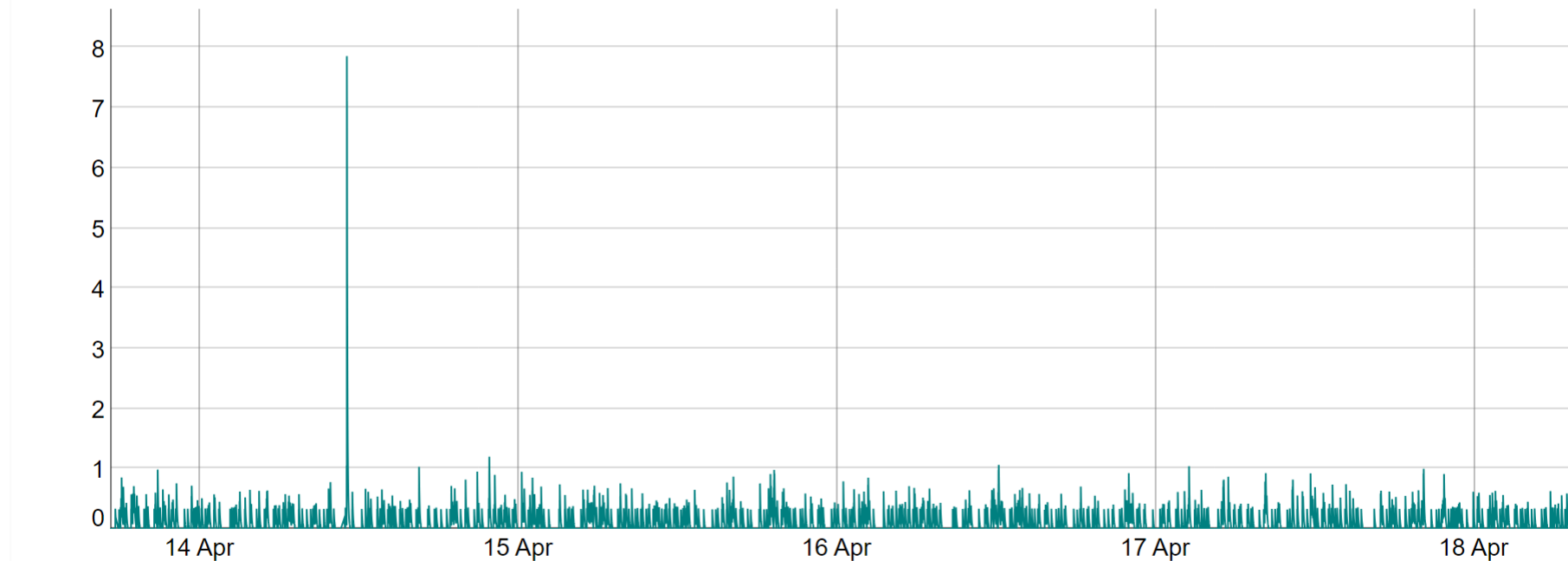
[data](#) | [query](#)

# Example: neutron burst

[show editor](#) | [hide editor](#) | [clear cache and refresh](#)

## Neutrons lab

### Neutron counts



time: not selected

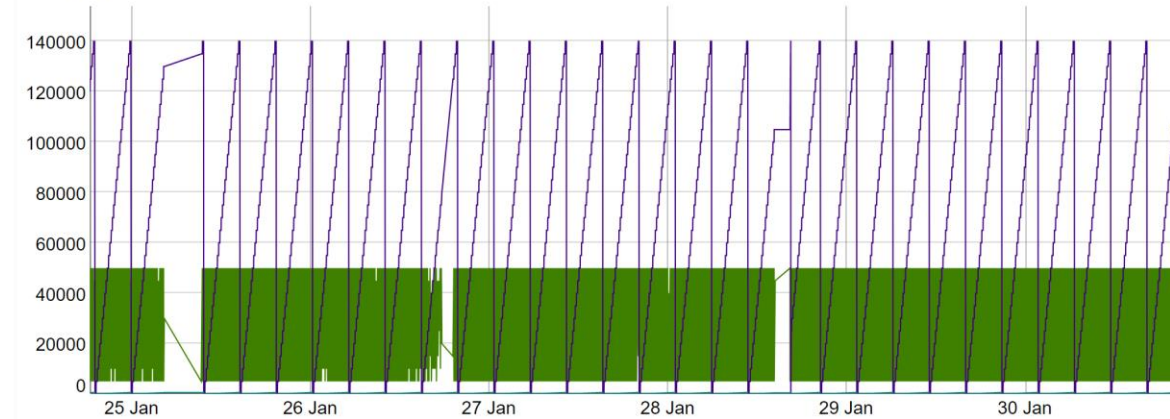
Neutrons

[log scale](#) | [lin scale](#)

[data](#) | [query](#)

# Example: plotting data from different data source

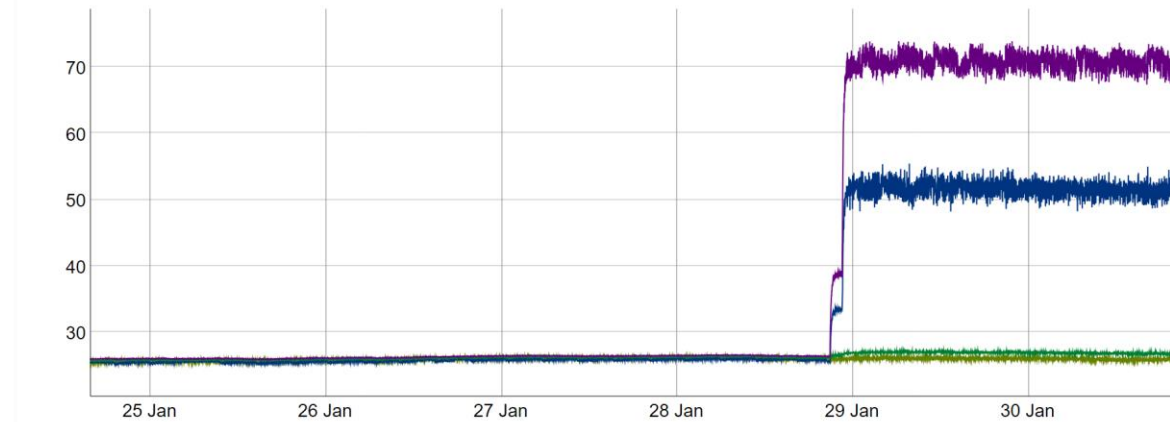
Stage position



time: not selected

- xvalue
  - zvalue
  - positioncounter
- [log scale](#) | [lin scale](#)  
[data](#) | [query](#)

Temperatures

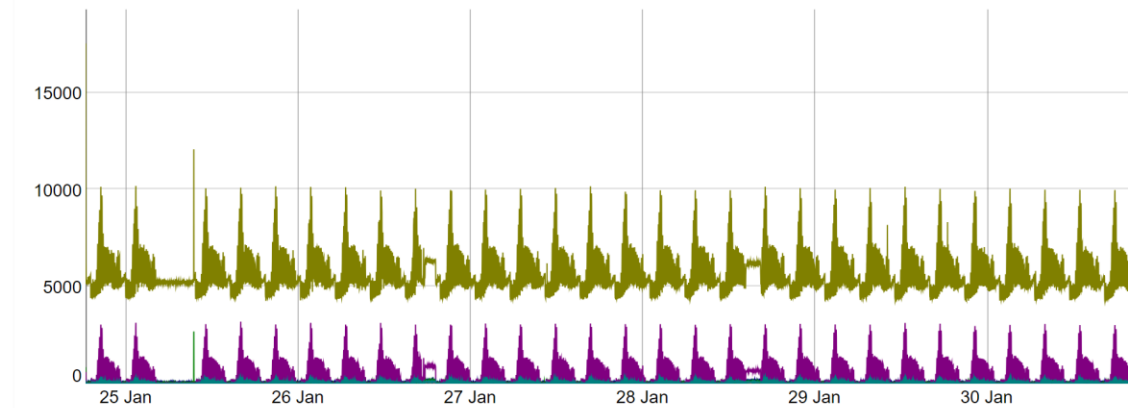


time: not selected

- Temp1
- Temp2

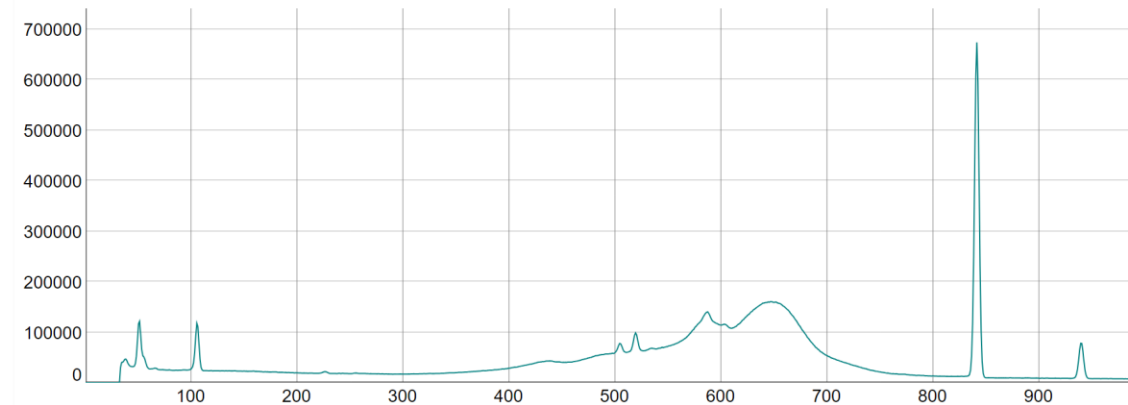
# Example: plotting data from different data source (cont'd)

Ortec time histories



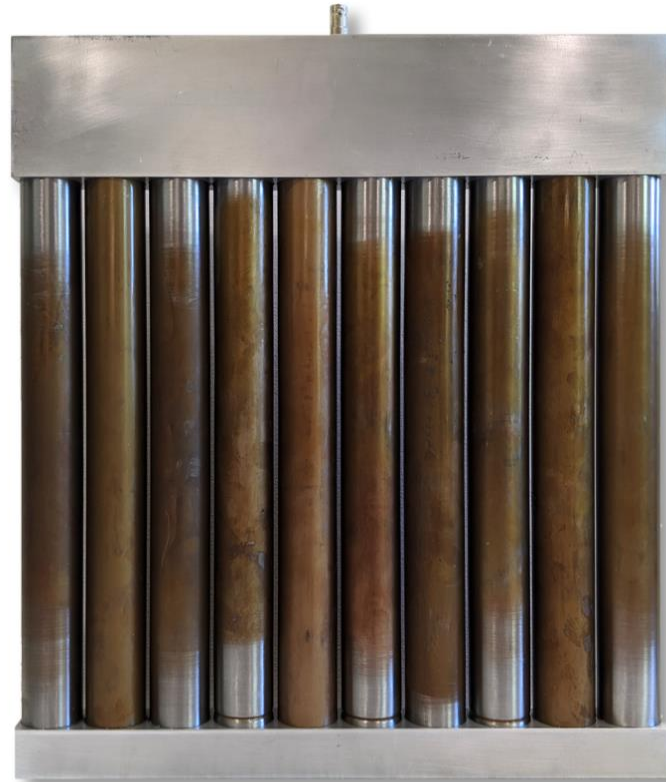
time: not selected  
 ch1-1100 (all)  
 ch90-120 (14 keV)  
 ch45-60 (6 keV)  
 ch835-850 (122 keV)  
 ch935-950 (136 keV)  
[log scale](#) | [lin scale](#)  
[data](#) | [query](#)

Ortec spectrum



# Obtaining low-cost large-area He-3 neutron detectors

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10x  $^3\text{He}$  Neutron Detector Bank

\$5,995.00

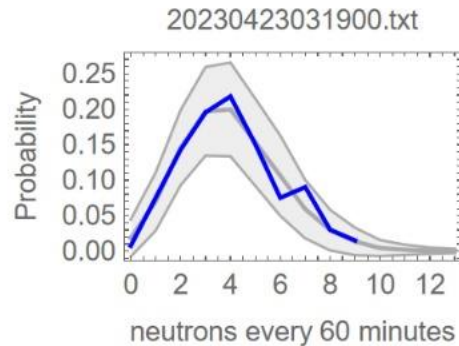


# Shielding to reduce background (neutrons, gammas)

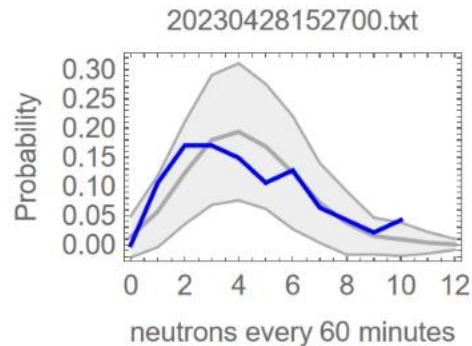
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- ▶ Working with Igor

# Toward a statistical framework for neutron/gamma detection



- Simulation of 100 experiments with 536 events and mean = 4.03008 (band == 2 sigma)
- Experiment with 536 events and mean 4.03008



- Simulation of 100 experiments with 200 events and mean = 4.25532 (band == 2 sigma)
- Experiment with 200 events and mean 4.25532

# Toward a statistical framework for isotopic analysis

Building on:

*Journal of Radioanalytical and Nuclear Chemistry, Vol. 263, No. 3 (2005) 691–696*

## Use of combined NAA and SIMS analyses for impurity level isotope detection

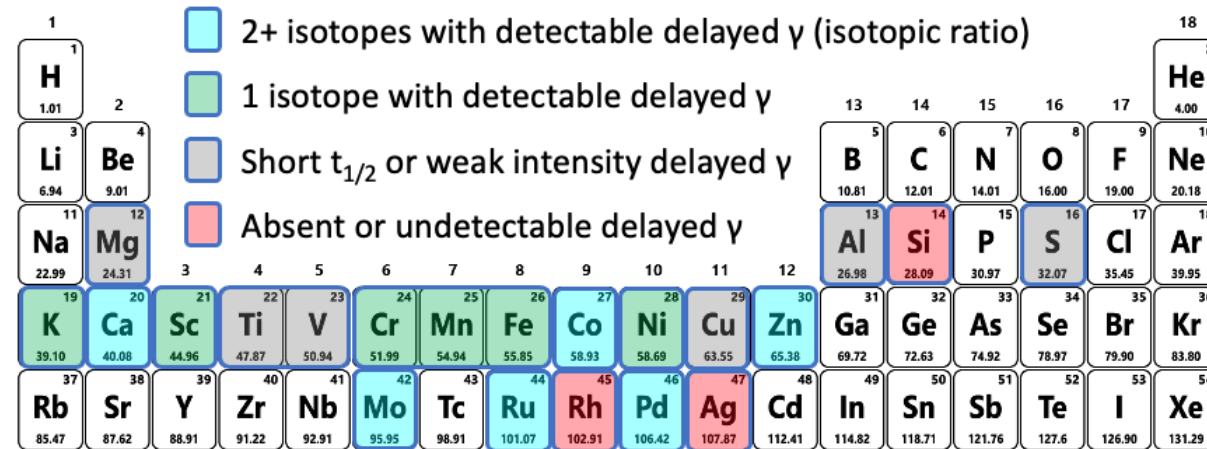
G. H. Miley,<sup>1\*</sup> G. Narne, T. Woo

*University of Illinois at Urbana-Champaign, Department of Nuclear, Plasma, and Radiological Engineering,  
103 S. Goodwin, Urbana, IL 61801, USA*

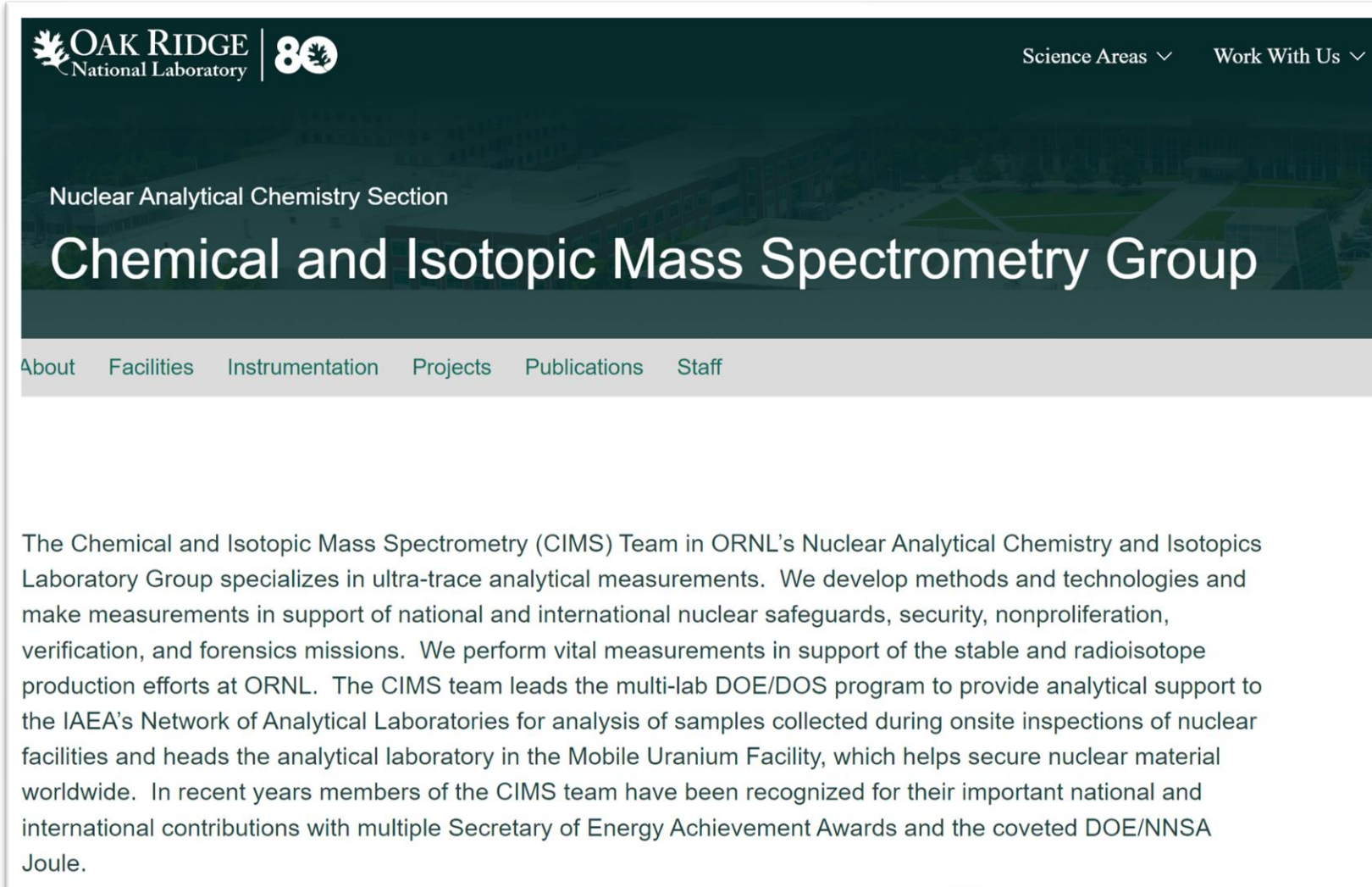
(Received April 6, 2004)

Neutron activation analysis (NAA) offers advantages for detecting impurity levels of select isotopes that have suitable neutron cross sections. Secondary ion mass spectrometry (SIMS) on the other hand detects most isotopes, but suffers various molecular interferences and covers only a small beam size volume per run. These two methods are combined here to study a large number of isotopes in titanium thin films in an electrolytic cell experiment. Nine isotopes are covered by NAA and over 50 with SIMS. An overlap in the data sets allows a normalization of SIMS data to the more accurate NAA measurements.

Promising isotopes:



# Toward a statistical framework for isotopic analysis (cont'd)



The screenshot shows the website for the Chemical and Isotopic Mass Spectrometry Group at ORNL. The header includes the ORNL logo and a 80th anniversary emblem. Navigation links for 'Science Areas' and 'Work With Us' are visible. The main content area features the group name and a list of navigation links: 'About', 'Facilities', 'Instrumentation', 'Projects', 'Publications', and 'Staff'. Below this is a detailed paragraph about the group's work.

OAK RIDGE National Laboratory | 80

Science Areas ▾ Work With Us ▾

Nuclear Analytical Chemistry Section

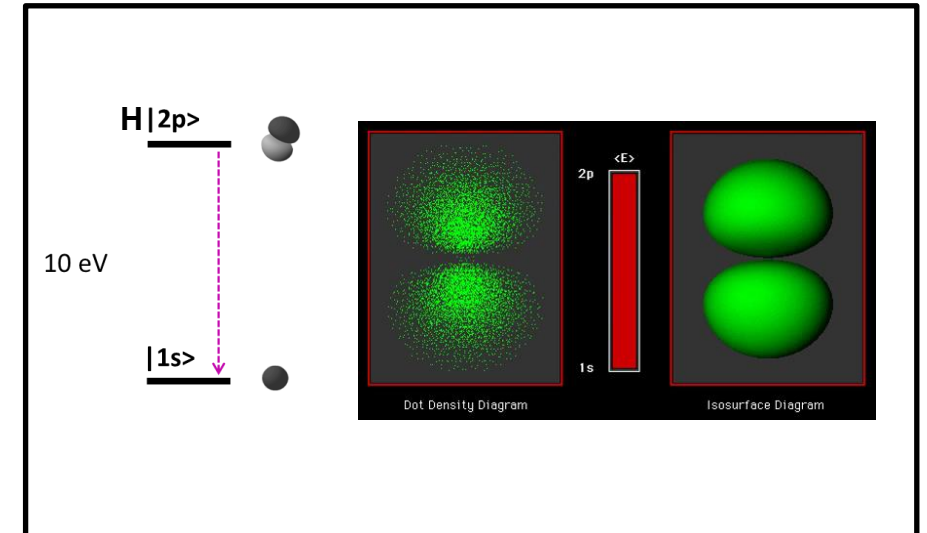
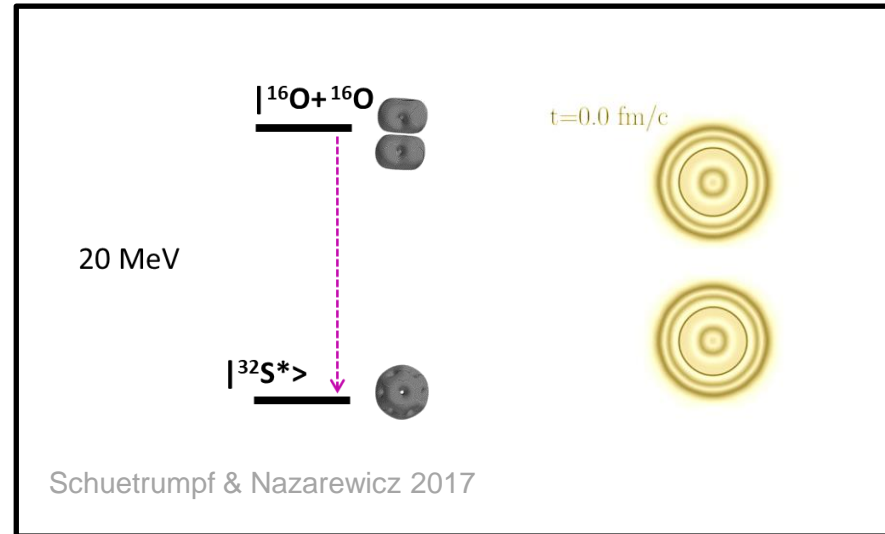
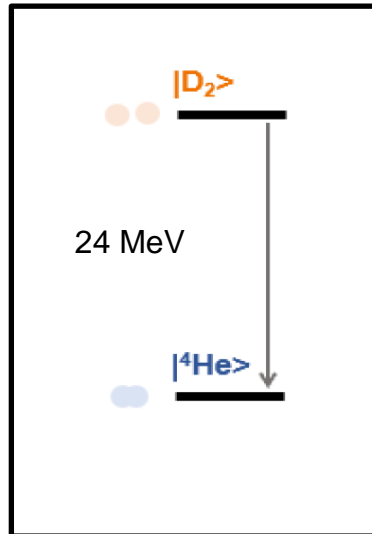
## Chemical and Isotopic Mass Spectrometry Group

[About](#) [Facilities](#) [Instrumentation](#) [Projects](#) [Publications](#) [Staff](#)

The Chemical and Isotopic Mass Spectrometry (CIMS) Team in ORNL's Nuclear Analytical Chemistry and Isotopics Laboratory Group specializes in ultra-trace analytical measurements. We develop methods and technologies and make measurements in support of national and international nuclear safeguards, security, nonproliferation, verification, and forensics missions. We perform vital measurements in support of the stable and radioisotope production efforts at ORNL. The CIMS team leads the multi-lab DOE/DOS program to provide analytical support to the IAEA's Network of Analytical Laboratories for analysis of samples collected during onsite inspections of nuclear facilities and heads the analytical laboratory in the Mobile Uranium Facility, which helps secure nuclear material worldwide. In recent years members of the CIMS team have been recognized for their important national and international contributions with multiple Secretary of Energy Achievement Awards and the coveted DOE/NNSA Joule.

# Framing

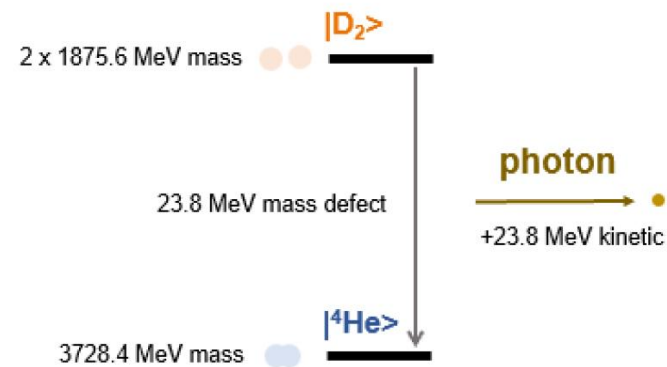
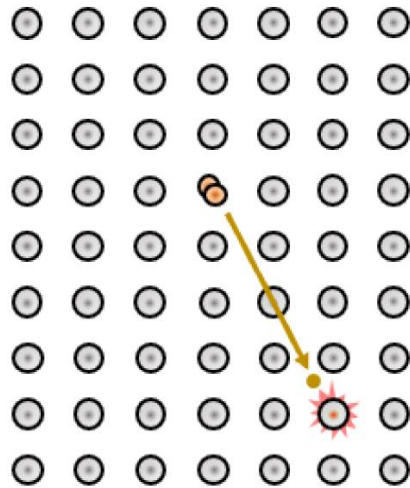
Fusion reaction viewed as a state transition to a lower energetic state, hindered by the reconfiguration of nucleons.



Analogous to how a state transition at the atomic level is hindered by the reconfiguration of electrons.

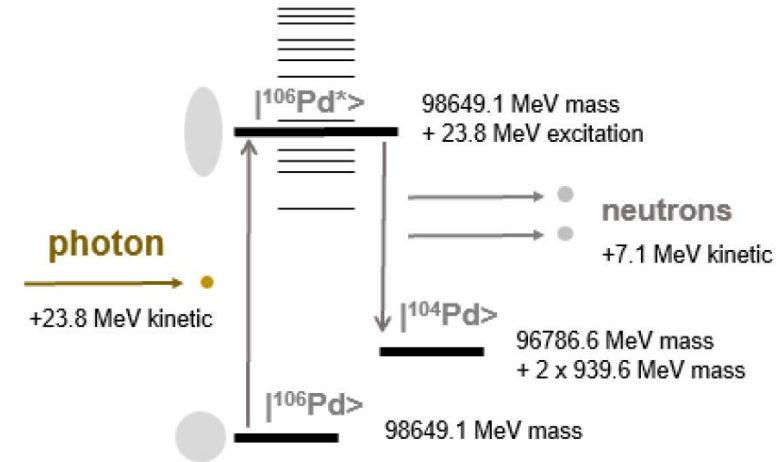
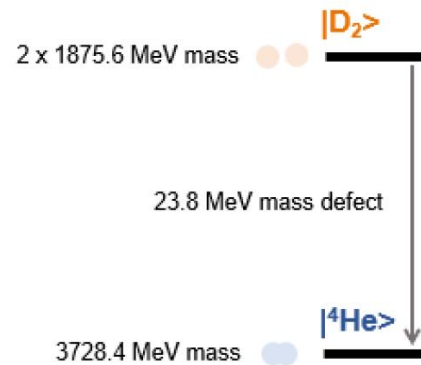
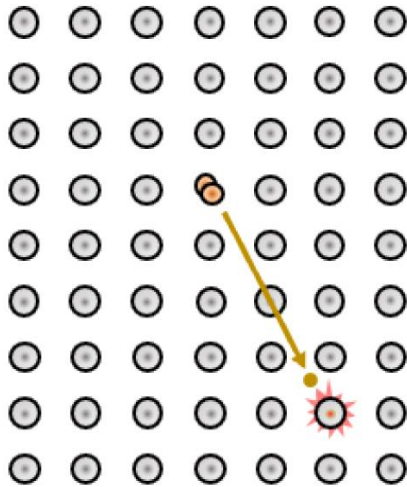
# Fusion-fission hybrid reactions

Conventionally expected but hopelessly slow:



# Fusion-fission hybrid reactions

Conventionally expected but hopelessly slow:



# Terhune & Baldwin PRL 1965

## NUCLEAR SUPERRADIANCE IN SOLIDS\*

J. H. Terhune and G. C. Baldwin

Advanced Technology Laboratories, General Electric Company, Schenectady, New York  
(Received 16 February 1965)

A general theory of coherent spontaneous gamma-ray emission from an assemblage of isomeric nuclei in a perfect crystalline solid has been developed. The solid, characterized by internal energy states of the nuclei, by the lattice vibrations, and by the electromagnetic field, is treated as an integrated quantized system rather than as a number of noninteracting nuclei.<sup>1-4</sup> Transition probabilities are calculated by the usual methods of first-order time-dependent perturbation theory.

Coherent spontaneous emission of radiation from a gas has been discussed by Dicke.<sup>5</sup> It was shown that transitions exist for which the radiation rates, line shapes, and linewidths are all different from the corresponding quantities for an assemblage of noninteracting radiators. In particular, certain states were predicted that possess radiation rates much greater than normal because of correlations among the internal motions of the various molecules composing the system.

In a solid composed of  $N$  identical two-level nuclei in a perfect crystal lattice at a uniform and low temperature, correlations in the internal motions of the radiators are more probable than in the case of a gas. Furthermore, the interactions among members of the solid system are much stronger than in the gas, because of the coupling between neighbors in the lattice. The usual assumption<sup>1-4</sup> that each nucleus radiates independently of the states of other nuclei in the system is incompatible with the coupling of the nuclei through the common electromagnetic and phonon fields. Calculations

of the spontaneous radiation rate for a solid system in which the nuclei are a priori assumed independent preclude the possibility of coherent spontaneous gamma emission by assumption. The present analysis is free from this inconsistency. Finally, the wavelength of the radiation is comparable with the spacing of nuclei in the lattice.

Using the method of Dicke,<sup>5</sup> the nuclear states are described by a vector model in which the vector orientation is quantized in energy space in analogy with fermion spin. The nuclei are assumed identical, in a uniform and field-free environment, with only two nondegenerate internal energy states coupled by a radiative transition. The lattice is assumed harmonic with nearest-neighbor interactions only; the phonon spectrum is approximated by the Debye model. The crystal is considered in the adiabatic approximation, and is assumed to be at rest with respect to the observer.

In the Hamiltonian for this system,

$$H = H_{\text{nuclei}} + H_{\text{lattice}} + H_{\text{radiation}} + H',$$

all terms except the interaction term  $H'$  are independent of the time. The latter is

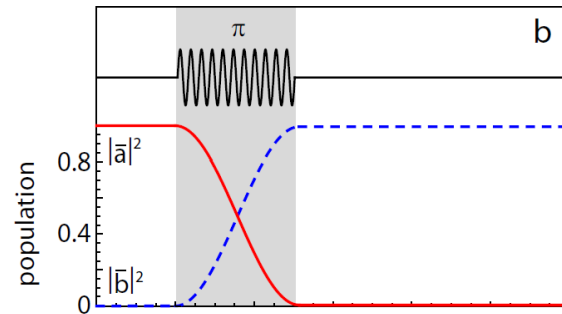
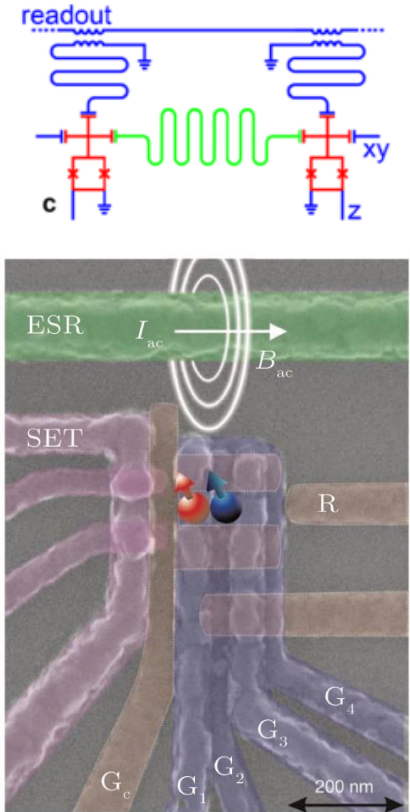
$$H' = -\frac{1}{2} \sum_k [(\vec{a}_k \cdot \vec{e})R_{k+} + (\vec{a}_k^* \cdot \vec{e}^*)R_{k-}],$$

in which  $\vec{a}_k^*$  and  $\vec{a}_k$  are photon creation and destruction operators, respectively,  $k$  characterizes a mode of the electromagnetic field,  $\vec{e}$  and  $\vec{e}^*$  were defined in reference 5, and the nuclear excitation and de-excitation operators



# Excitation transfer

## Quantum state transfer through activation of couplings

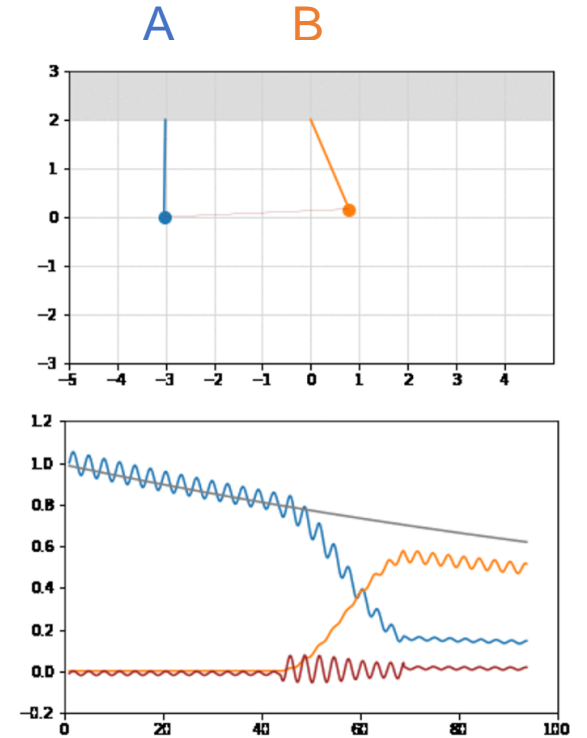


Frimmer & Novotny (2014)  
*Am. Journal of Physics*

Two qubits  
implemented in silicon.

Zhang et al. (2018) *Chinese Physics B*

kinetic energy (classical)  $\triangleq$   
state occupation probability  
(quantum)



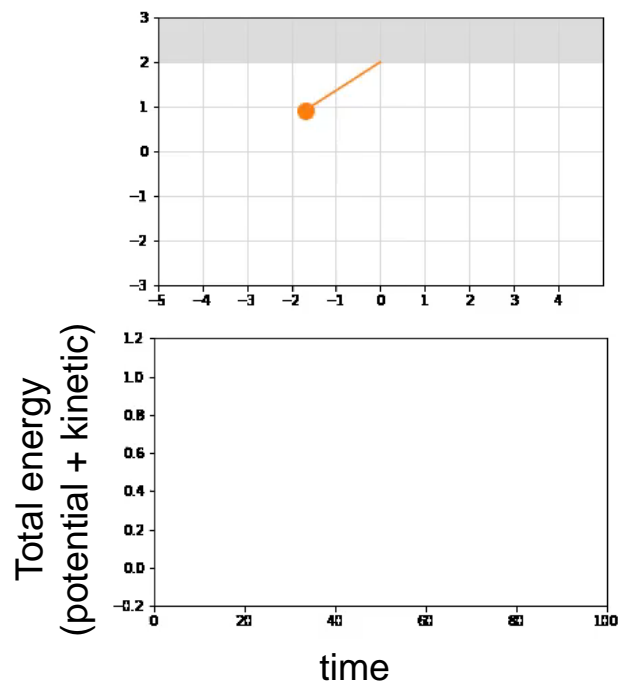
Classical analog to two  
coupled quantum systems

Briggs et al. (2011) *Physical Review E*  
Eisfeld et al. (2012) *Physical Review E*

# Mechanical analogs for excitation transfer

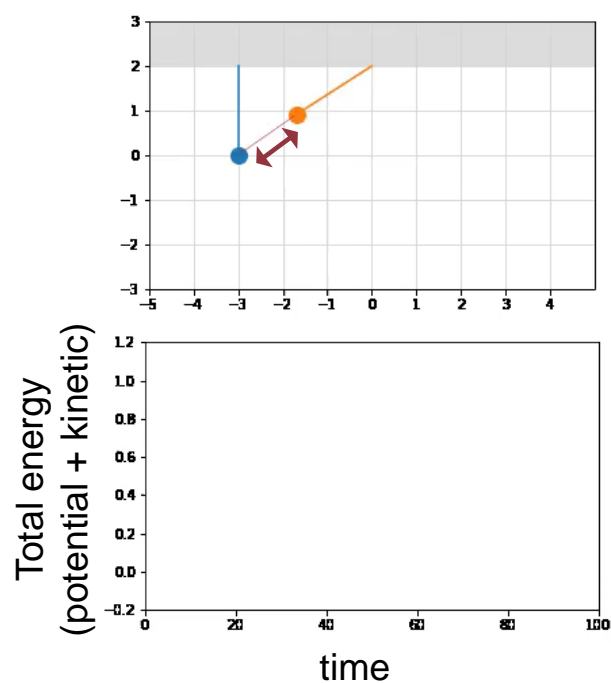
Damped oscillations

Start: **Orange** excited



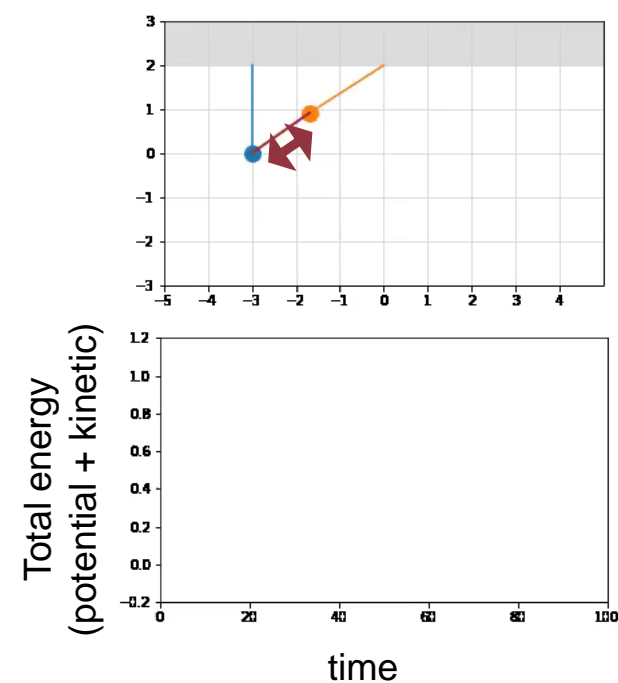
Weak coupling  $\leftrightarrow$

Start: **Blue** non-excited



Strong coupling  $\leftrightarrow$

Start: **Blue** non-excited

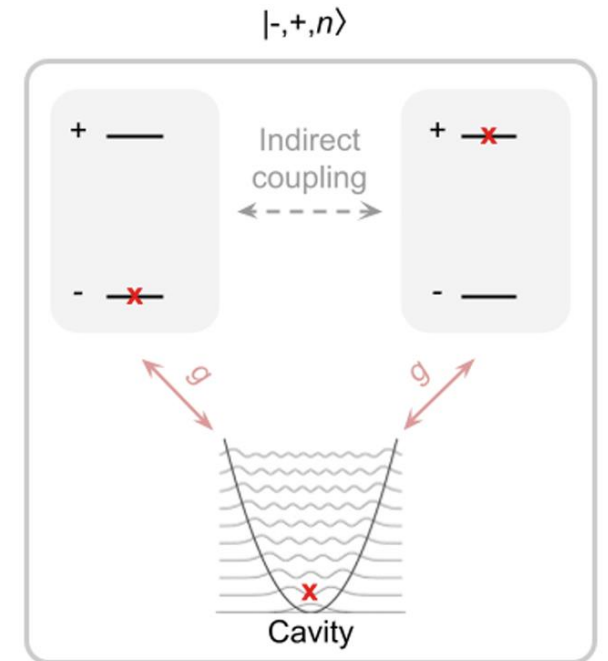
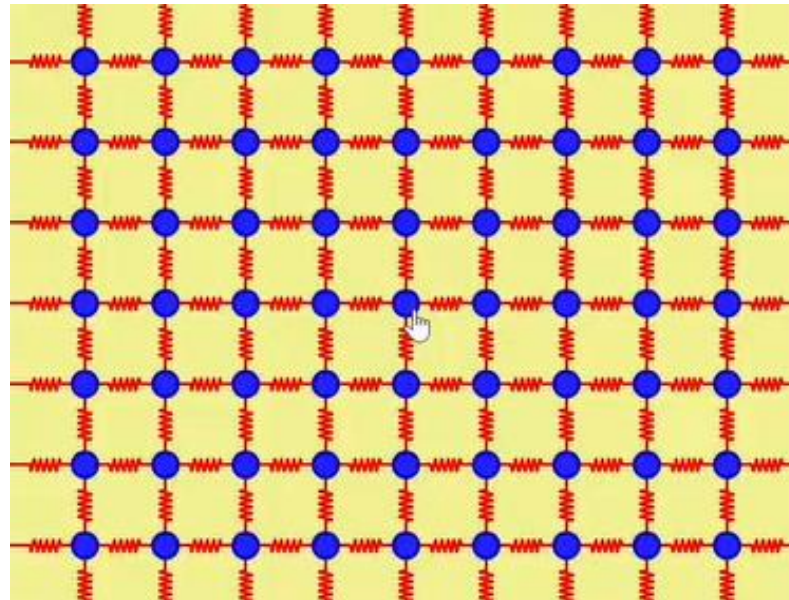
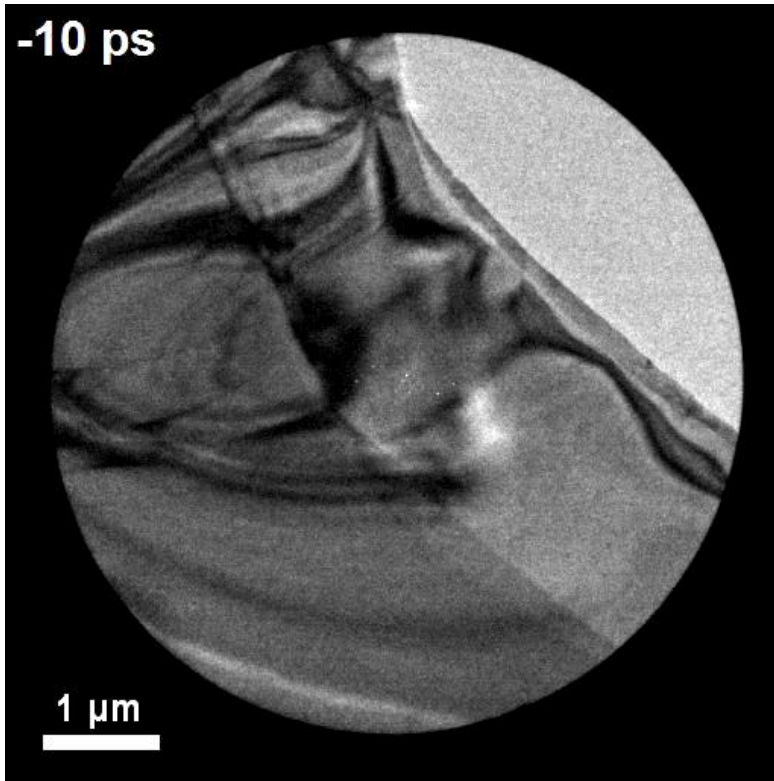


- Energy (total)
- Energy of orange pendulum
- Energy of blue pendulum
- Energy in the elastic string

Briggs and Einfeld, *PRA* (2013)  
 Zimanyia and Silbey, *J. Chem. Phys.* (2010)

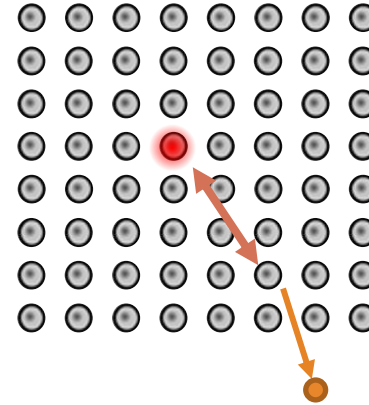
# Phonons/plasmons as source of couplings

Phonons in Ge



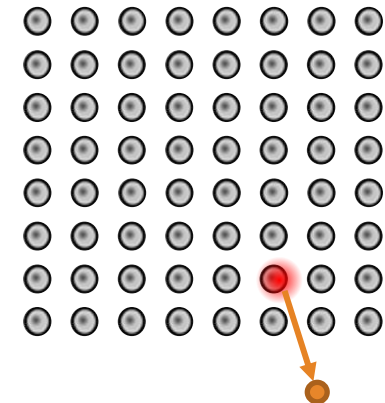
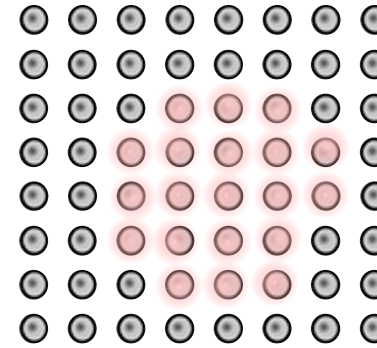
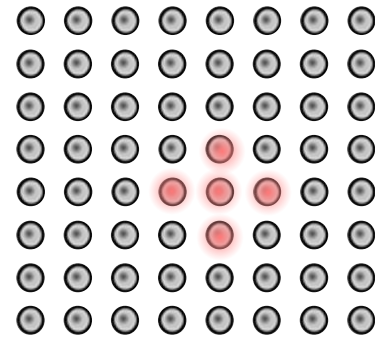
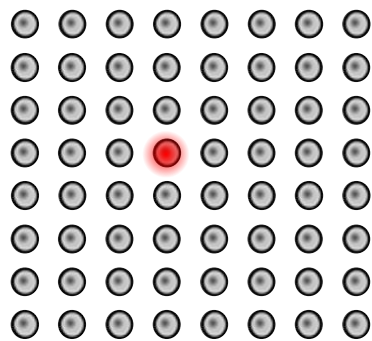
Cremons et al. 2016, *Nat. Commun.*

# Delocalized excited states as intermediate steps



(one possible pathway of many)

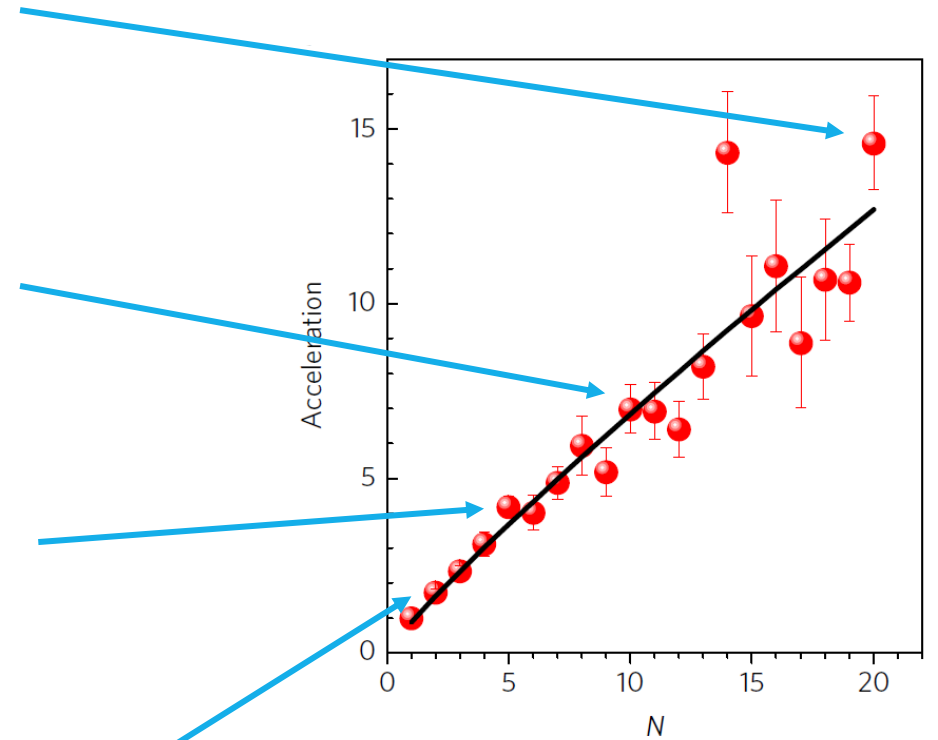
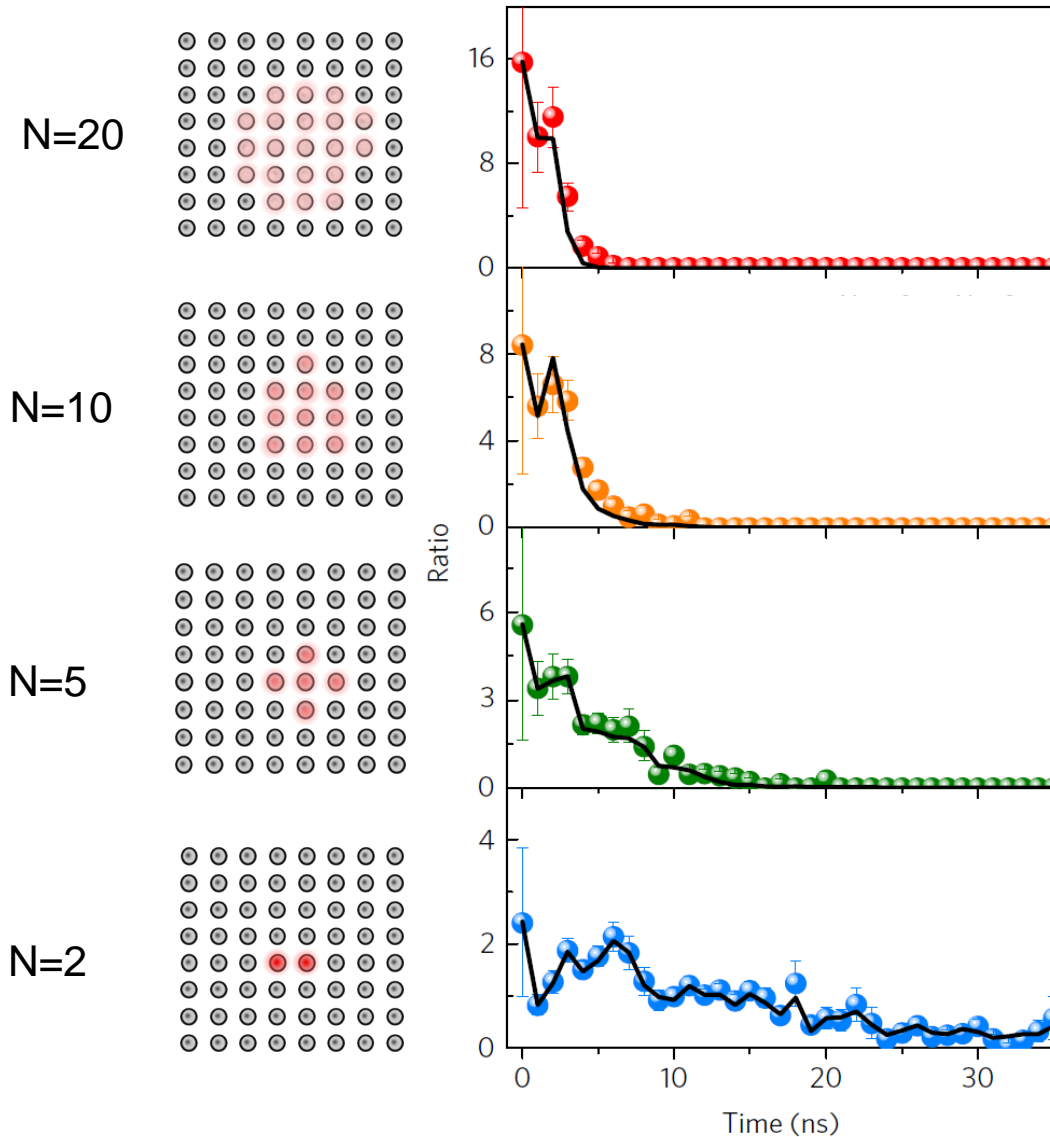
$$|\uparrow\rangle = \frac{1}{\sqrt{N}}(|\uparrow\downarrow\downarrow\dots\downarrow\rangle + |\downarrow\uparrow\downarrow\dots\downarrow\rangle + \dots + |\downarrow\downarrow\downarrow\dots\uparrow\rangle)$$



coupling activated

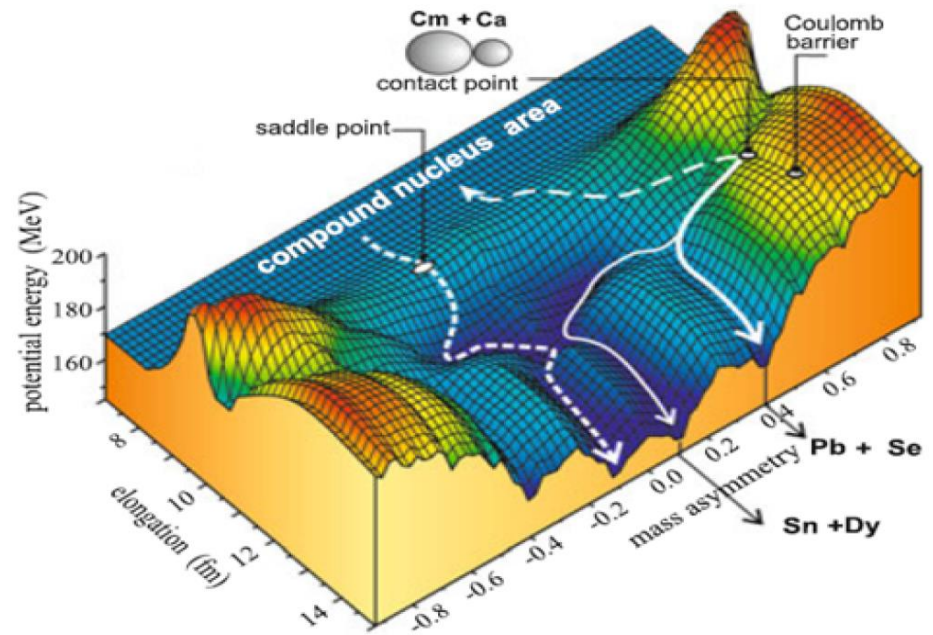
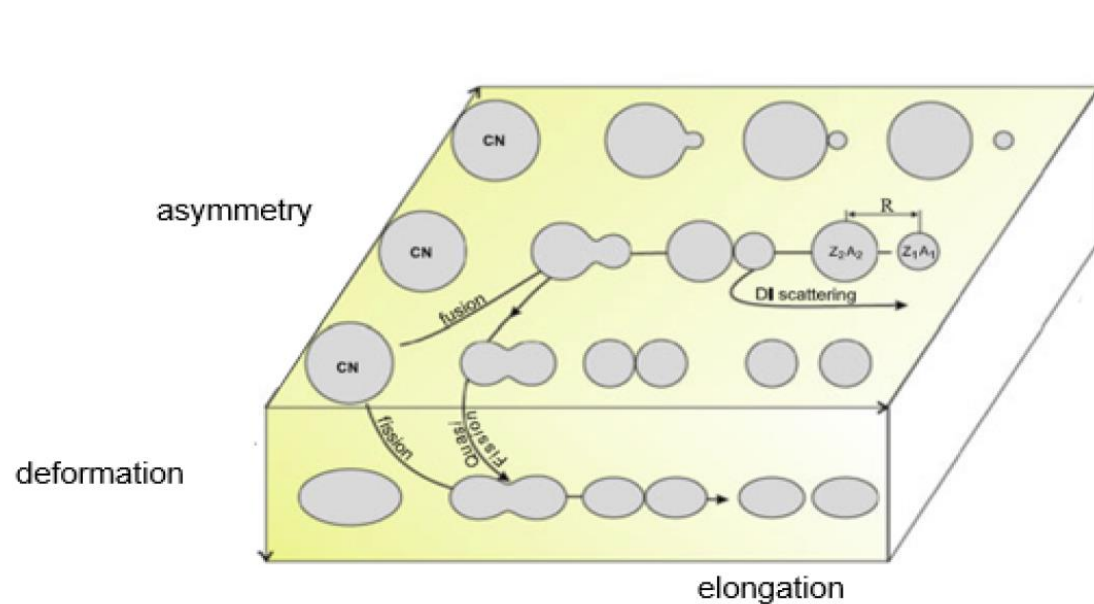
# Superradiance: coherent acceleration of emission

Couplings associated with delocalized states



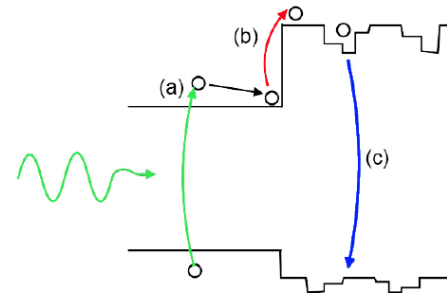
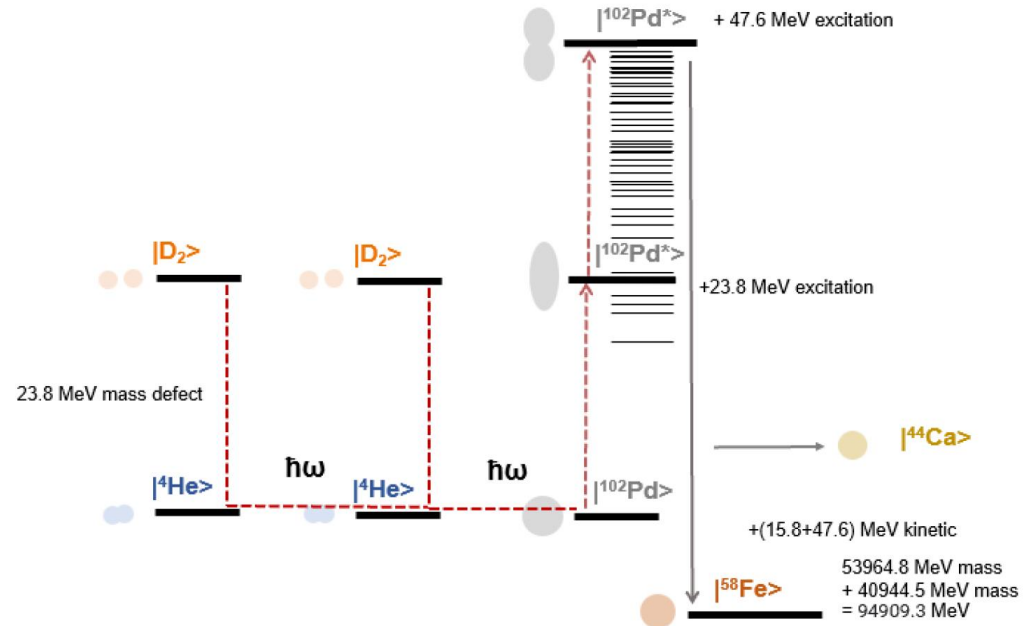
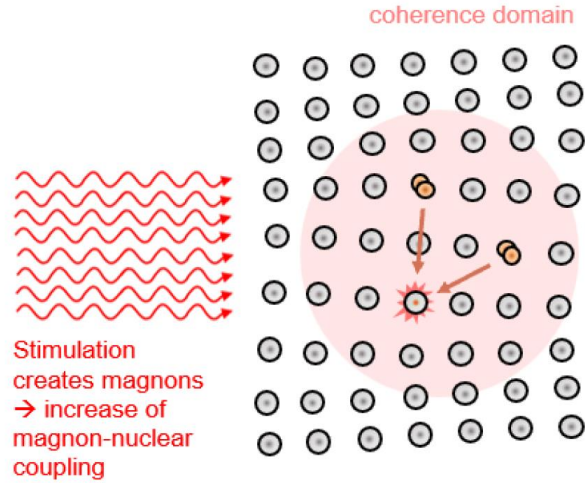
# What determines nuclear energy levels & reaction products

Nuclear molecule literature:



Zagrebaev & Greiner 2010, Springer

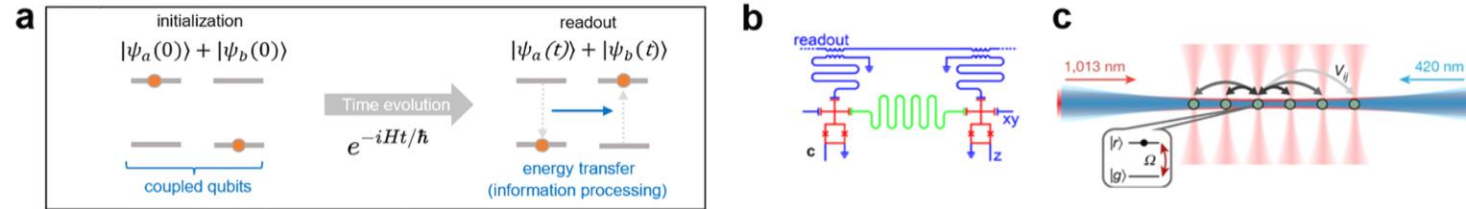
# Low-Z elements from near-symmetric Pd fission



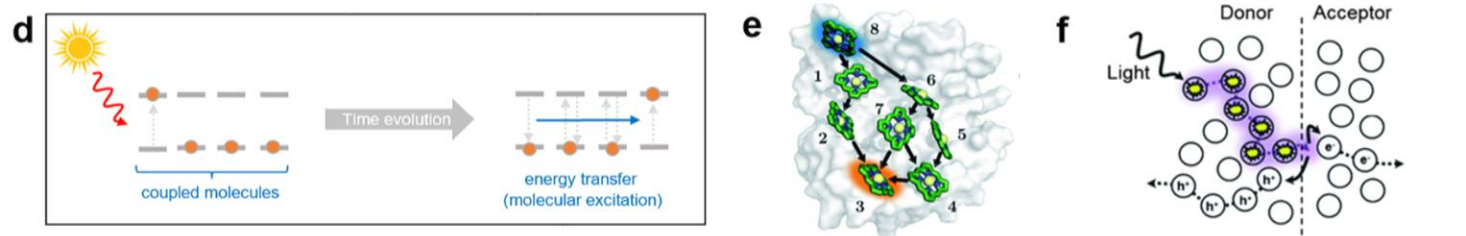
Quantum ratcheting

# The Emergence of Quantum Energy Science

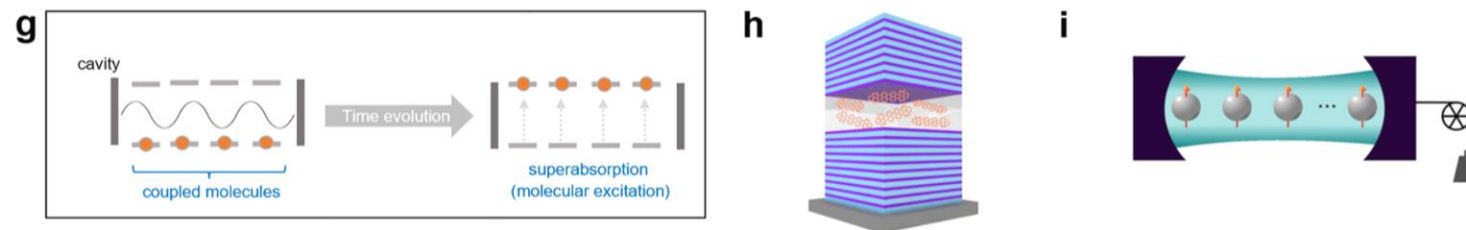
## Quantum computing



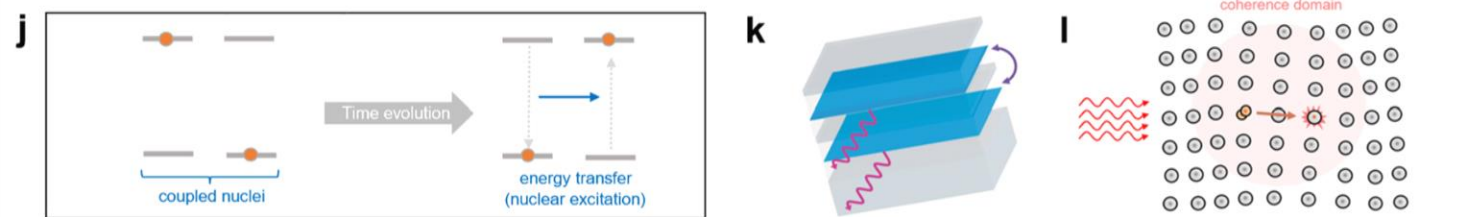
## Quantum solar



## Quantum batteries




## Quantum nuclear



Metzler et al. 2022  
<https://arxiv.org/abs/2303.01632>




# First conference on "Quantum Energy"




**International Conference on Quantum Energy**  
4-6 December 2023  
Pullman on the Park  
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86 days 17 hrs 9 mins 3 secs  
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**Dr James Q. Quach** | PhD, PGradDip, BEng (Hons), BCS, BCom  
Science Leader in Quantum Science and Technologies  
Australia's national science agency, CSIRO

SCIENCE ADVANCES | RESEARCH ARTICLE

### PHYSICS

## Superabsorption in an organic microcavity: Toward a quantum battery

**James Q. Quach<sup>1\*</sup>, Kirsty E. McGhee<sup>2</sup>, Lucia Ganzer<sup>3</sup>, Dominic M. Rouse<sup>4</sup>, Brendon W. Lovett<sup>4</sup>, Erik M. Gauger<sup>5</sup>, Jonathan Keeling<sup>4</sup>, Giulio Cerullo<sup>3</sup>, David G. Lidzey<sup>2</sup>, Tersilla Virgili<sup>3\*</sup>**

The rate at which matter emits or absorbs light can be modified by its environment, as markedly exemplified by the widely studied phenomenon of superradiance. The reverse process, superabsorption, is harder to demonstrate because of the challenges of probing ultrafast processes and has only been seen for small numbers of atoms. Its central idea—superextensive scaling of absorption, meaning larger systems absorb faster—is also the key idea underpinning quantum batteries. Here, we implement experimentally a paradigmatic model of a quantum battery, constructed of a microcavity enclosing a molecular dye. Ultrafast optical spectroscopy allows us to observe charging dynamics at femtosecond resolution to demonstrate superextensive charging rates and storage capacity, in agreement with our theoretical modeling. We find that decoherence plays an important role in stabilizing energy storage. Our work opens future opportunities for harnessing collective effects in light-matter coupling for nanoscale energy capture, storage, and transport technologies.

Conference dates

# Accelerating chemical reactions via quantum effects

*“Scientists are intensely interested in what are known as ‘quantum-enhanced’ chemical reactions”*

Phys.org 2023

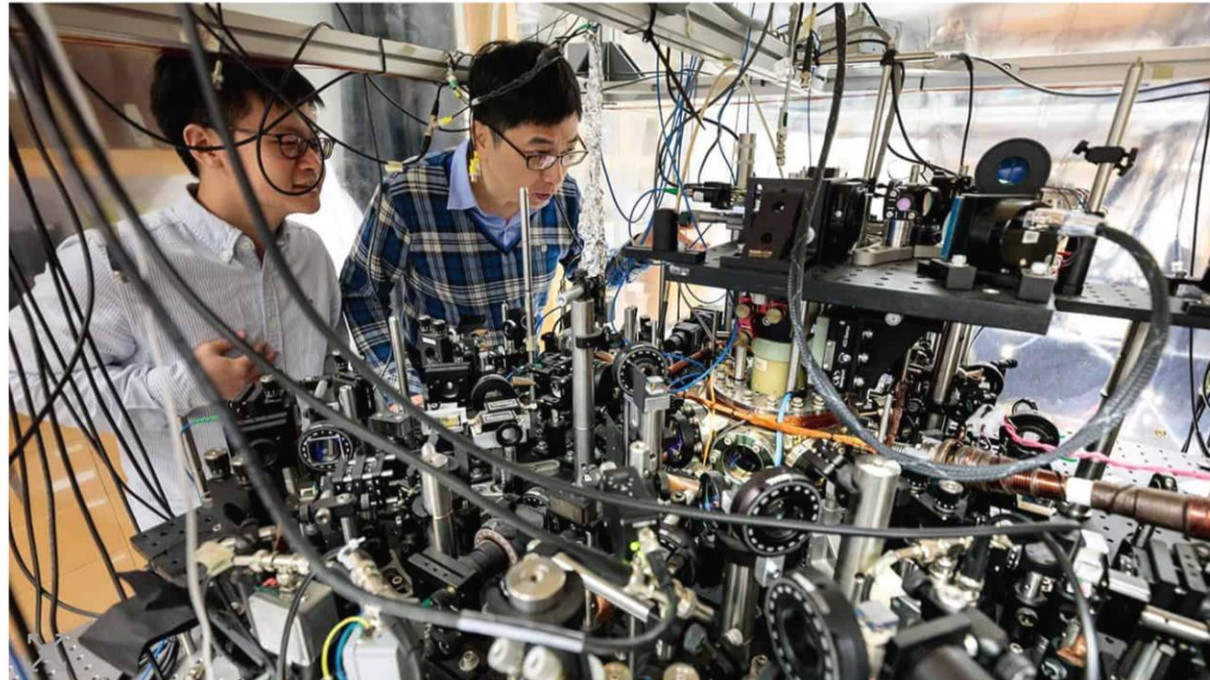
*“molecules sharing a quantum state might produce accelerated chemical reactions if those molecules were ‘coupled’ together and reacting as one”*

sciencealert 2023

CHEMICAL PROCESSES | RESEARCH UPDATE

## Quantum superchemistry emerges in the laboratory

23 Aug 2023

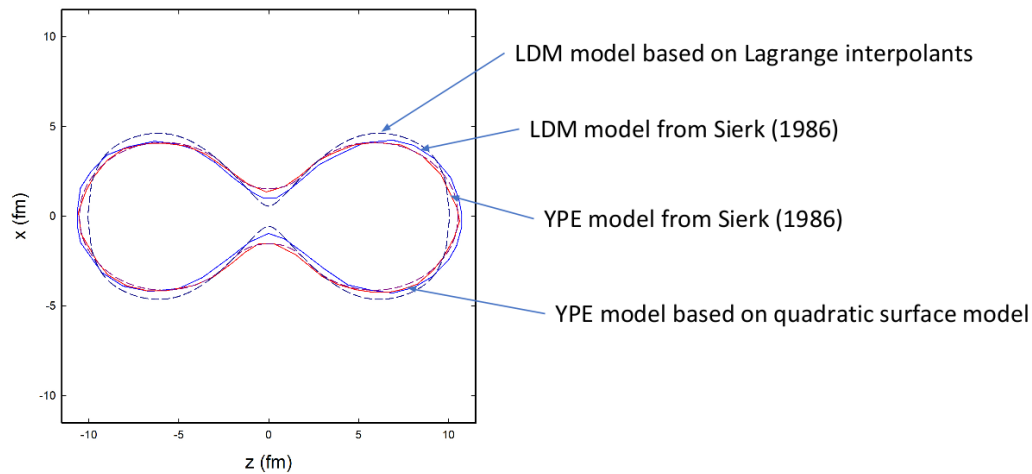


Atoms and molecules: Cheng Chin (r) and postdoctoral researcher Zhendong Zhang in the University of Chicago laboratory where they and colleagues observed the first evidence of quantum superchemistry. (Courtesy: John Zich/University of Chicago)

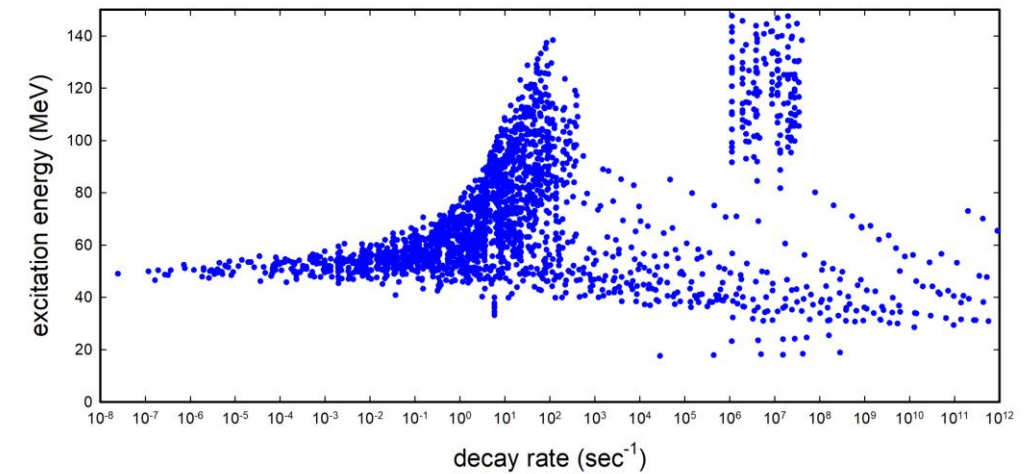
# Liquid drop models for determining nuclear excited states

## Pd-106 nuclear molecule (Vd-53 + Vd-53)

- Bohr-Wheeler model predicts a (not very stable) symmetric nuclear molecule for Pd-106 near 40.7 MeV



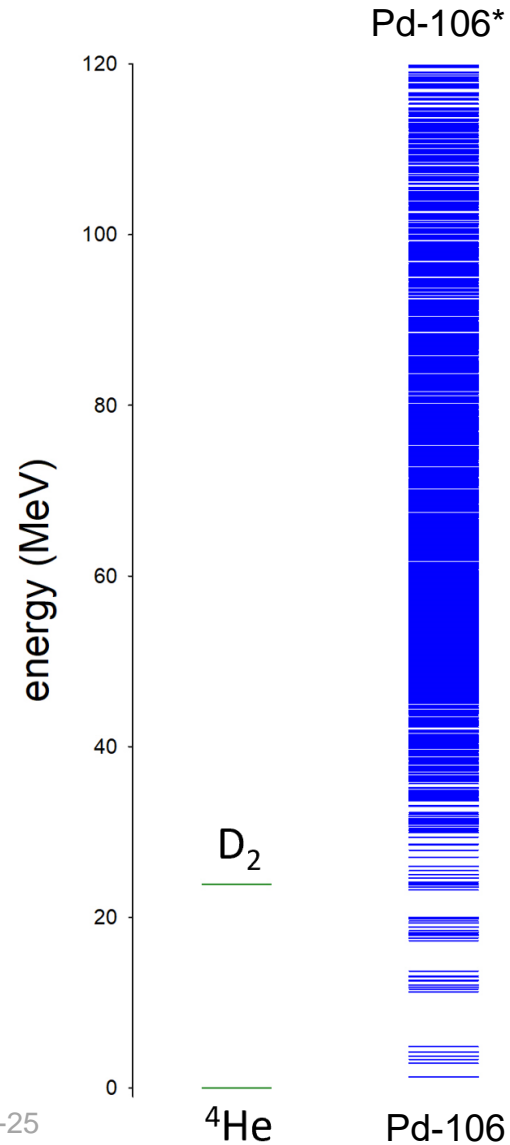
## Nuclear molecule energy vs decay rate including tunneling and daughter decays



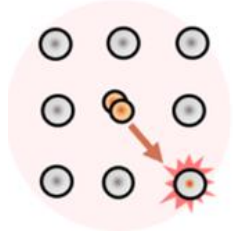
For all the stable Pd isotopes

Hagelstein 2023, ICCF-25

# Long-lived excited states in Pd available for excitation transfer



# Intermediate variables



Hamiltonian: 
$$H = \underbrace{\sum_j M_j c^2}_{\text{nuclei}} + \underbrace{\hbar \omega a^\dagger a}_{\text{osc}} + \underbrace{\sum_j (-\mu_j \cdot B)_+ + (-\mu_j \cdot B)_-}_{\text{magn. interaction}}$$

$10^{12}$   $D_2$  molecules in PdD vacancies  
( $U_\theta = 150$  eV, fluctuations  $\pm 5$  pm)  
coupled to  $10^6$  Pd nuclei via magnon-nuclear  
coupling (coupling strength  $V = 100$  neV)

# QuTIP Python modeling of excitation transfer

Jupyter  
nbviewer

JUPYTER FAQ </> [Menu] [Refresh] [Close] [Download]

```
In [37]: bra_labels, ket_labels = make_braket_labels(nmm_list)
```

Now let's plot the results using a helper function `plot_prob` that pulls together plotting code that we used in the last tutorial.

```
In [38]: plot_prob(P, times, ket_labels)
plt.title(f"2 TSS with {H_latex} ( $\Delta E \approx 1.9328$ ,  $\omega=1$ ,  $U=0.1$ ) (Fig 8)");
```

2 TSS with  $H = \Delta E/2(\sigma_{z1} + \sigma_{z2}) + \hbar\omega(a^\dagger a + 1/2) + U(a^\dagger + a)(\sigma_{x1} + \sigma_{x2})$  ( $\Delta E \approx 1.9328$ ,  $\omega = 1$ ,  $U = 0.1$ ) (Fig 8)

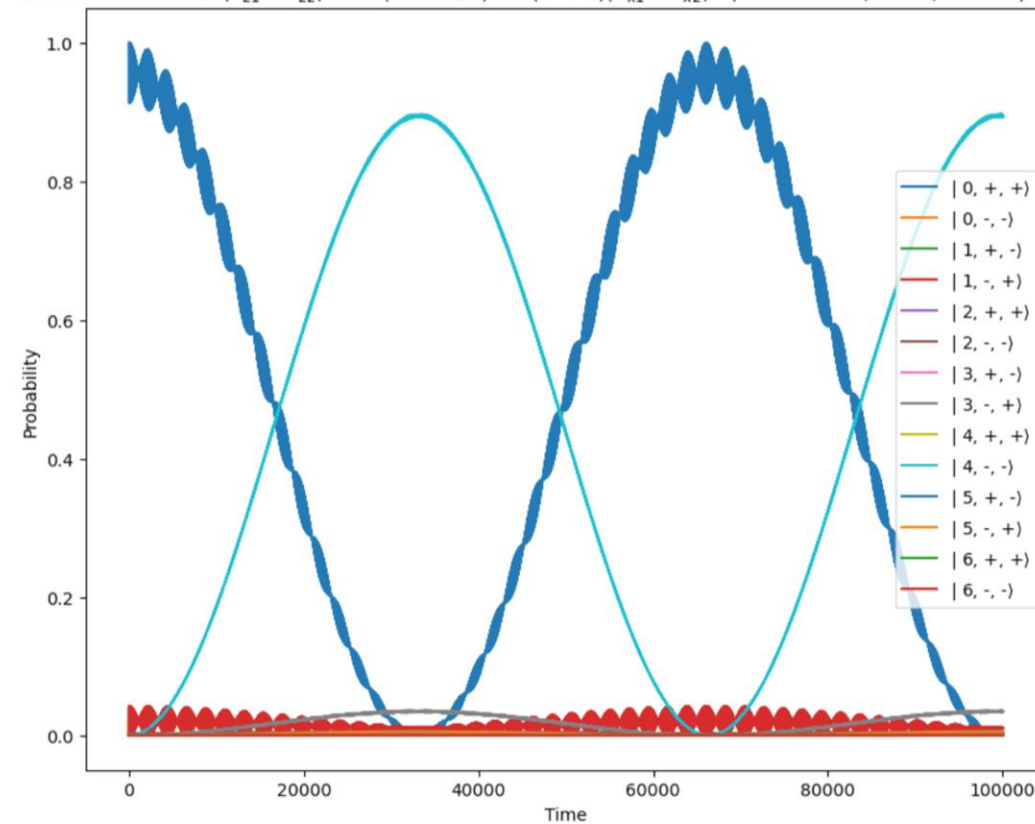
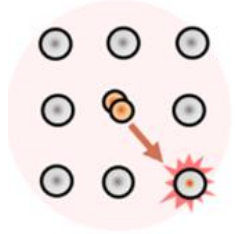


Fig 8 shows exactly what we expected, namely down conversion - both TSS transition from "+" to "-" each giving of 2 bosons in the process i.e.  $|0, +, +\rangle \rightarrow |4, -, -\rangle$ .

<https://github.com/project-ida/two-state-quantum-systems/>

# Intermediate variables



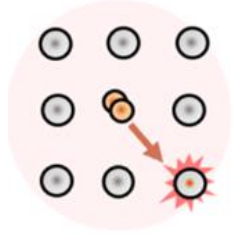
**Dependent variable:**

Hamiltonian: 
$$H = \underbrace{\sum_j M_j c^2}_{\text{nuclei}} + \underbrace{\hbar\omega a^\dagger a}_{\text{osc}} + \underbrace{\sum_j (-\mu_j \cdot B)_+ + (-\mu_j \cdot B)_-}_{\text{magn. interaction}}$$

$10^{12}$   $D_2$  molecules in PdD vacancies  
 ( $U_e = 150$  eV, fluctuations  $\pm 5$  pm)  
 coupled to  $10^6$  Pd nuclei via magnon-nuclear  
 coupling (coupling strength  $V = 100$  neV)

Transfer rate = fusion rate: 
$$\Gamma = \sqrt{N_{D_2}} \sqrt{N_{Pd-106}} \frac{[\frac{V}{\Delta E} \sqrt{\frac{vol_{nuc}}{vol_{mol}}} e^{-G}] \frac{V}{\Delta E} \Phi_{Pd-106}}{\hbar} = \frac{10^{-7} eV 10^{-17} 10^{-6} 10^{-7} eV}{24 \times 10^6 eV} 1 \underbrace{\sqrt{10^{12}} \sqrt{10^6} \frac{1}{\hbar}}_{\text{Dicke enhancement}} = 10^{-18} s^{-1}$$

# Intermediate variables



Hamiltonian: 
$$H = \underbrace{\sum_j M_j c^2}_{\text{nuclei}} + \underbrace{\hbar\omega a^\dagger a}_{\text{osc}} + \underbrace{\sum_j (-\mu_j \cdot B)_+ + (-\mu_j \cdot B)_-}_{\text{magn. interaction}}$$

$10^{12}$  D<sub>2</sub> molecules in PdD vacancies  
 ( $U_\theta = 150$  eV, fluctuations +/- 5pm)  
 coupled to  $10^6$  Pd nuclei via magnon-nuclear  
 coupling (coupling strength  $V = 100$  neV)

**Dependent variable:**

Transfer rate  
 = fusion rate:

$$\Gamma = \sqrt{N_{D_2}} \sqrt{N_{Pd-106}} \frac{[\frac{V}{\Delta E} \sqrt{\frac{vol_{nuc}}{vol_{mol}}} e^{-G}] \frac{V}{\Delta E} \Phi_{Pd-106}}{\hbar} = \frac{10^{-7} eV 10^{-17} 10^{-6} 10^{-7} eV}{24 \times 10^6 eV} 1 \underbrace{\sqrt{10^{12}} \sqrt{10^6} \frac{1}{\hbar}}_{\text{Dicke enhancement}} = 10^{-18} s^{-1}$$

**Intermediate variables:**

# of nuclei  
 coupled to  
 oscillator

Energy in  
 oscillator  
 mode

DD proximity &  
 screening potential

Closeness of resonance between  
 D<sub>2</sub> (23.8 MeV) and lattice nuclei

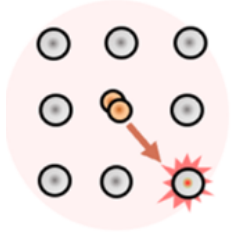
Excited plasmon  
 & phonon modes

D diffusion  
 rate

Lattice structure  
 (vacancy content)



# Intermediate variables



Hamiltonian: 
$$H = \underbrace{\sum_j M_j c^2}_{\text{nuclei}} + \underbrace{\hbar\omega a^\dagger a}_{\text{osc}} + \underbrace{\sum_j (-\mu_j \cdot B)_+ + (-\mu_j \cdot B)_-}_{\text{magn. interaction}}$$

$10^{12}$  D<sub>2</sub> molecules in PdD vacancies  
 ( $U_e = 150$  eV, fluctuations +/- 5pm)  
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 coupling (coupling strength  $V = 100$  neV)

**Dependent variable:**

Transfer rate  
 = fusion rate:

$$\Gamma = \sqrt{N_{D_2}} \sqrt{N_{Pd-106}} \frac{[\frac{V}{\Delta E} \sqrt{\frac{vol_{nuc}}{vol_{mol}}} e^{-G}] \frac{V}{\Delta E} \Phi_{Pd-106}}{\hbar} = \frac{10^{-7} eV 10^{-17} 10^{-6} 10^{-7} eV}{24 \times 10^6 eV} \underbrace{1 \sqrt{10^{12}} \sqrt{10^6} \frac{1}{\hbar}}_{\text{Dicke enhancement}} = 10^{-18} s^{-1}$$

**Intermediate variables:**

# of nuclei coupled to oscillator  
 Energy in oscillator mode  
 DD proximity & screening potential  
 Closeness of resonance between D<sub>2</sub> (23.8 MeV) and lattice nuclei

Excited plasmon & phonon modes  
 D diffusion rate  
 Lattice structure (vacancy content)

**Independent variables:**

Laser focal point size  
 Laser intensity, wavelength, pulse length & freq.  
 Temperature & pressure  
 Microstructure of untreated sample  
 Isotopic composition of untreated sample

# All the building blocks are there and accepted

ARTICLES

<https://doi.org/10.1038/s41566-017-0013-3>

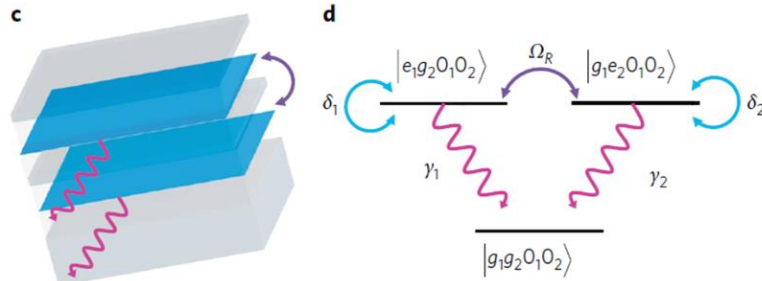
nature  
photonics

## Rabi oscillations of X-ray radiation between two nuclear ensembles

Johann Haber<sup>1</sup>, Xiangjin Kong<sup>2,3</sup>, Cornelius Strohm<sup>1</sup>, Svenja Willing<sup>1</sup>, Jakob Gollwitzer<sup>1</sup>, Lars Bocklage<sup>1</sup>, Rudolf Ruffer<sup>4</sup>, Adriana Pálffy<sup>2\*</sup> and Ralf Röhlsberger<sup>1\*</sup>

The realization of the strong coupling regime between a single cavity mode and an electromagnetic resonance is a centrepiece of quantum optics. In this regime, the reversible exchange of a photon between the two components of the system leads to so-called Rabi oscillations. Strong coupling is used in the optical and infrared regimes, for instance, to produce non-classical states

Rabi oscillations of the Rabi oscillations spectrum. Our results



Article

## Coherent X-ray–optical control of nuclear excitons

<https://doi.org/10.1038/s41586-021-03276-x>

Received: 22 October 2018

Accepted: 2 December 2020

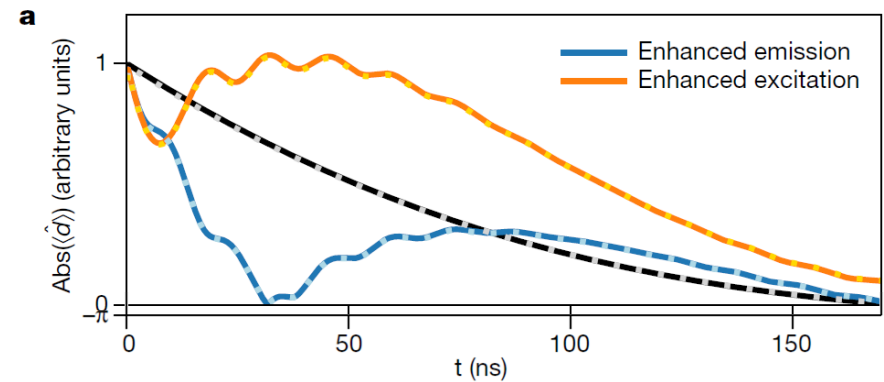
Published online: 17 February 2021

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Kilian P. Heeg<sup>1</sup>, Andreas Kaldun<sup>1</sup>, Cornelius Strohm<sup>2</sup>, Christian Ott<sup>1</sup>, Rajagopalan Subramanian<sup>1</sup>, Dominik Lentrodt<sup>1</sup>, Johann Haber<sup>2</sup>, Hans-Christian Wille<sup>2</sup>, Stephan Goettler<sup>1</sup>, Rudolf Ruffer<sup>2</sup>, Christoph H. Keitel<sup>1</sup>, Ralf Röhlsberger<sup>2,4,5,6,7</sup>, Thomas Pfeifer<sup>8</sup> & Jörg Evers<sup>1,9\*</sup>

Coherent control of quantum dynamics is key to a multitude of fundamental studies and applications<sup>1</sup>. In the visible or longer-wavelength domains, near-resonant light fields have become the primary tool with which to control electron dynamics<sup>2</sup>. Recently, coherent control in the extreme-ultraviolet range was demonstrated<sup>3</sup>, with a few-attosecond temporal resolution of the phase control. At hard-X-ray energies (above 5–10 kiloelectronvolts), Mössbauer nuclei feature narrow nuclear resonances due to their recoilless absorption and emission of light, and spectroscopy of these resonances is widely used to study the magnetic, structural and dynamical properties of matter<sup>4,5</sup>. It



# All the building blocks are there and accepted

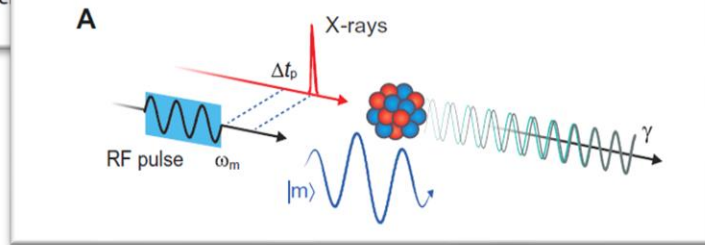
SCIENCE ADVANCES | RESEARCH ARTICLE

PHYSICS

## Coherent control of collective nuclear quantum states via transient magnons

Lars Bocklage<sup>1,2\*</sup>, Jakob Gollwitzer<sup>1</sup>, Cornelius Strohm<sup>1</sup>, Christian F. Adloff<sup>1,2</sup>, Kai Schlage<sup>1</sup>, Ilya Sergeev<sup>1</sup>, Olaf Leupold<sup>1</sup>, Hans-Christian Wille<sup>1</sup>, Guido Meier<sup>3,2</sup>, Ralf Röhlsberger<sup>1,2,4,5,6</sup>

Ultrafast and precise control of quantum systems at x-ray energies involves photons with oscillation periods below 1 as. Coherent dynamic control of quantum systems at these energies is one of the major challenges in hard x-ray quantum optics. Here, we demonstrate that the phase of a quantum system embedded in a solid can be coherently controlled via a quasi-particle with subattosecond accuracy. In particular, we tune the quantum phase of a collectively excited nuclear state via transient magnons with a precision of 1 zs and a timing stability below 50 fs. These small temporal shifts are monitored interferometrically via quantum beats between different hyperfine-split levels. The experiment demonstrates zeptosecond interferometry and shows that transient quasi-particles e

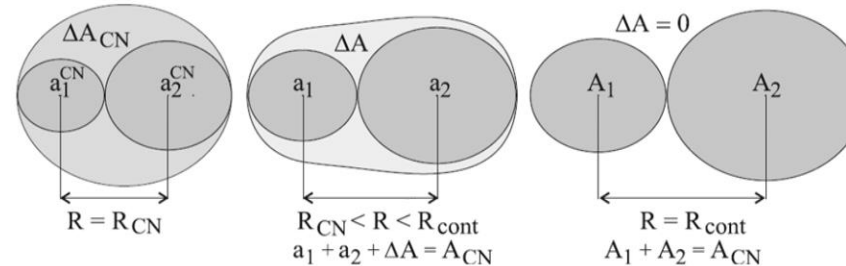


## Chapter 7 Giant Nuclear Systems of Molecular Type

Valery Zagrebaev and Walter Greiner

### 7.1 Introduction

Cluster structure is very often set off against the shell structure of light and medium nuclei. However the appearance of clusters themselves (compact pieces of nuclear matter) is conditioned just by the shell effects. In light nuclei these clusters are mainly alpha-particles. In heavy nuclear systems tightly packed nuclei (such as <sup>132</sup>Sn or <sup>208</sup>Pb) may lead to energetically favorable two (and even three) center configurations. These cluster configurations play an important role both in the structure of heavy nuclear systems and in the low-energy nuclear dynamics. The asymmetric nuclear fission (see, for example, Ref. [1]), the heavy-ion radioactivity [2, 3], the shape isomeric states of heavy nuclei [4] and the true ternary fission of superheavy nuclei (see below) are the manifestations of such kind of clusterization. Our studies of fusion–fission reactions and multi-nucleon transfer processes in low-energy heavy ion collisions demonstrated that the shell



# All the building blocks are there and accepted

nature  
physics

LETTERS

<https://doi.org/10.1038/s41567-017-0001-z>

## Superradiance of an ensemble of nuclei excited by a free electron laser

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In 1954 Dicke predicted the accelerated initial decay of multiple atomic excitations<sup>1</sup>, laying the foundation for the concept of superradiance. Further studies<sup>2–4</sup> suggested that emission of the total energy was similarly accelerated, provided that the system reaches the inversion threshold. Superradiant emission of the total energy has been confirmed by numerous studies<sup>4–12</sup>, yet the acceleration of the initial decay has not been experimentally demonstrated. Here we use resonant diffraction of X-rays from the Mössbauer transition<sup>13</sup> of <sup>57</sup>Fe nuclei to investigate superradiant decay, photon by photon, along the entire chain of the de-excitation cascade of up to 68 simultaneous coherent nuclear excitations created by a pulse of an X-ray free-electron laser. We find agreement with Dicke's theory<sup>1</sup> for the accelerated initial decay as the number of excitations is increased. We also find that our results are in agreement with a simple statistical model, providing a necessary baseline for discussing further properties of superradiance, within and beyond the low-excitation regime.

Even when  $N \ll n_e$ , this is a very large  $N$ -fold acceleration. Therefore, the acceleration of the initial decay predicted by Dicke can be studied even in the low-excitation regime.

We studied the accelerated initial decay of multiple coherent nuclear excitations created by an X-ray pulse of the SPring-8 Angstrom Compact free electron LASer (SACLA)<sup>14</sup>, the only source that can presently provide temporally and spatially coherent pulses of many photons within the bandwidth of the convenient 14.4 keV nuclear transition of <sup>57</sup>Fe. For X-rays, the small-system limit is fundamentally excluded because the wavelength is similar to inter-atomic distances. However, one can create a phased excitation of an extended system. The ideal 'X-ray lattice'<sup>10</sup> is offered by atomic periodicity, and the 'seeded coherence'<sup>12</sup> is provided in nuclear resonance diffraction conditions<sup>15</sup>. Similar to the 'collective dipole' of atoms coupled to the light field in an infrared optical cavity<sup>10–12</sup>, the 'compound' excited state<sup>16,17</sup> of an ensemble of nuclei under diffraction<sup>15</sup> or forward-scattering<sup>18,19</sup> conditions leads to enhancement of emission and strong 'speed-up' of the collective response<sup>15–23</sup>. In

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### Symmetry-enhanced supertransfer of delocalized quantum states

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**Abstract.** Coherent hopping of excitation relies on quantum coherence over physically extended states. In this work, we consider simple models to examine the effect of symmetries of delocalized multi-excitation states on the dynamical timescales, including hopping rates, radiative decay and environmental interactions. While the decoherence (pure dephasing) rate of an extended state over  $N$  sites is comparable to that of a non-extended state, superradiance leads to a factor of  $N$  enhancement in decay and absorption rates. In addition to superradiance, we illustrate how the multi-excitonic states exhibit 'supertransfer' in the far-field regime—hopping from a symmetrized state over  $N$  sites to a symmetrized state over  $M$  sites at a rate proportional to  $MN$ . We argue that such symmetries could play an operational role in physical systems based on the competition between symmetry-enhanced interactions and localized